

C9 Nanoscale Magnetism

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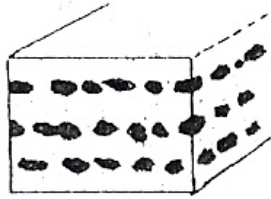
Introduction

Magnetic nanostructures have at least one dimension in the 1 - 100 nm size range. Magnetic properties differ from bulk

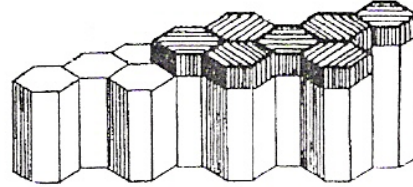
three small dimensions



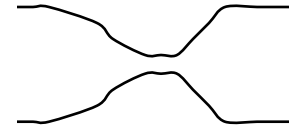
Nanoparticles



Nanocomposite

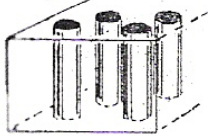


recording medium

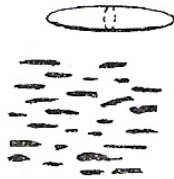


nanoconstriction

two small dimensions



Nanowires

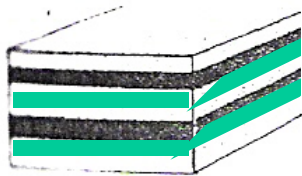


a

one small dimension



thin film



Thin film stack

-

Ratio of surface to volume atoms

Nanoparticle: $4\pi R^2 a / (4/3)\pi R^3 = 3a/R$

Nanowire: $2\pi R a / \pi R^2 = 2a/R$

Unsupported film: $2a/t$

If $2R, t \approx 10$ nm, $a \approx 0.25$ nm surface fraction is

➤ 15% for a nanoparticle

➤ 10 % for a nanowire

➤ 5% for a thin film

Characteristic length scales

Magnetic length scales are expressed in terms of the *exchange length* ($l_{ex} \approx 2 \text{ nm}$)

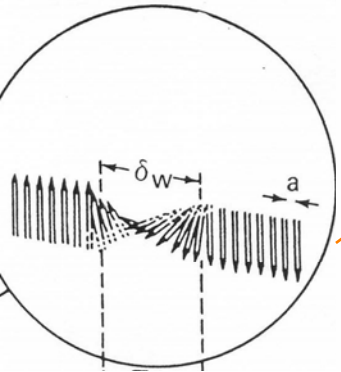
$$l_{ex} = \sqrt{(A/\mu_0 M_s^2)}$$

and the *hardness parameter* (>1 for a permanent magnet; $\ll 1$ for a soft magnet)

$$\kappa = \sqrt{(K/\mu_0 M_s^2)}$$

Table 8.1 Characteristic micromagnetic length scales (in nm).

Length	Expression	Fe ($\kappa = 0.12$)	Nd ₂ Fe ₁₄ B ($\kappa = 1.54$)
l_{ex}	$\sqrt{(A/\mu_0 M_s^2)}$	1.5	1.9
R_{coh}	$\sqrt{(24)l_{ex}}$	7	9
δ_B	$\pi l_{ex}/\kappa$	40	3.9
R_{sd}	$36\kappa l_{ex}$	6	107
R_{eq}	$(3k_B T/4\pi B M)$	0.8	0.9
R_b	$(6k_B T/K_1)^{1/3}$	8	2



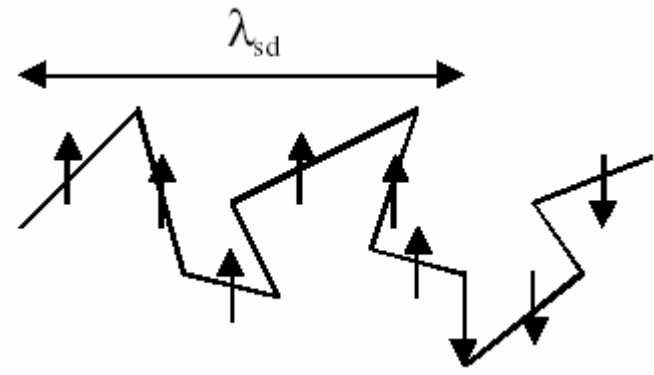
- l_{ex} exchange length
- R_{coh} maximum particle size for coherent rotation
- δ_B Bloch wall width
- R_{sd} maximum equilibrium single domain particle size
- R_{eq} particle radius for which $k_B T = MB$ in 1 tesla at 300 K
- R_b superparamagnetic blocking radius at 300 K.

- Transport lengths:

Mean free path $\lambda_{\uparrow} \lambda_{\downarrow}$ (1-10 nm, $\lambda_{\uparrow}/\lambda_{\downarrow}$ up to 5 in 3d transition metals)

Spin diffusion length $\lambda_{sd\uparrow} \lambda_{sd\downarrow}$

Typically an electron is scattered ~ 100 times before experiencing a spin flip, hence $\lambda_{sd} \approx 10\lambda$.



- Quantum length $l_B = \sqrt{\hbar/eB} = 26/\sqrt{B}$ nm (B in tesla)

Arises in quantum phenomena such as Landau diamagnetism

3 Superparamagnetism

Tiny particles (≈ 10 nm) are unstable when the barrier to magnetization reversal is comparable to $k_B T$

$$\Delta \rightarrow \Delta \pm \mu_0 m H \cos \theta$$

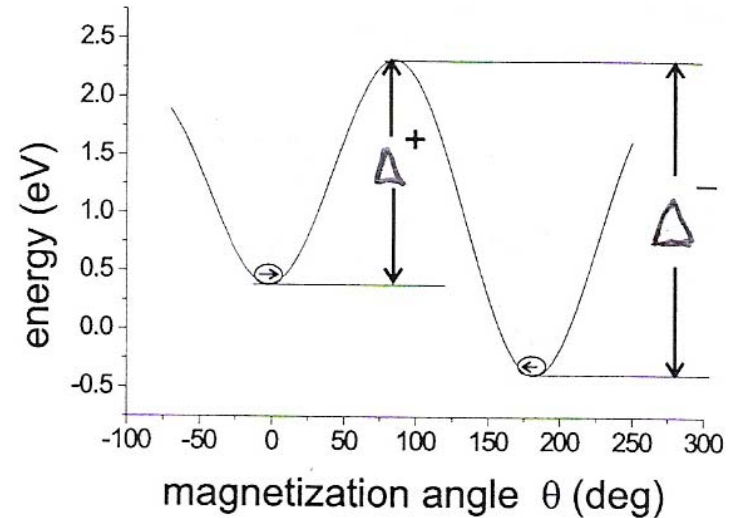
$$\tau = \tau_0 \exp (\Delta / k_B T)$$

$$\sim 10^{-10} \text{ s}$$

Δ May originate from magnetocrystalline anisotropy $K_1 V$, shape anisotropy $K_d V$, or surface anisotropy $K_s A$.

The magnetization decays exponentially,

$$M(t) = M(0) \exp(-t/\tau)$$

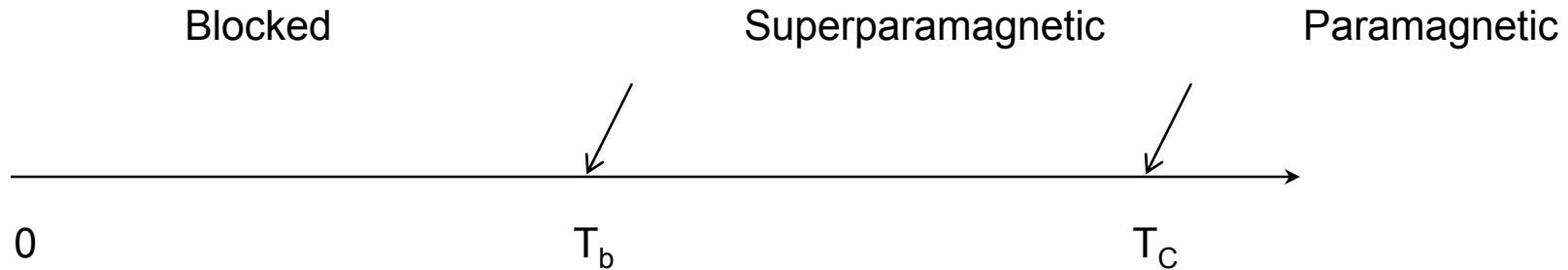


Radius	Relaxation time
3 nm	1.9 ms
4 nm	223 hr
5 nm	4.10^{12} y

or a cobalt particle of radius 3.5 nm at different temperatures

Temperature	relaxation time
260 K	332 s
300 K	10s
340 K	0.6s
380 K	76 ms

Blocking



Criteria for blocking $\Delta/k_B = 25 \Rightarrow \tau \approx 100\text{s}$

$\Delta/k_B = 40 \Rightarrow \tau \approx 10\text{y}$

In the superparamagnetic region the particles behave like a classical paramagnet with a giant classical moment m . The superparamagnetic susceptibility is

$$\chi = \mu_0 N m^2 / 3 k_B T$$

, where N is the number of particles per cubic meter.

On cooling basalt in the Earth's magnetic field H_e , tiny superparamagnetic particles of magnetite block below T_B , thereby acquiring a *thermoremanent magnetization*.

$$M_{tr} = \mu_0 N m^2 H / 3 k_B T_b$$

Polarity of earth's field changes randomly every $\approx 100,000$ y

Plates move at about 1 cm y^{-1} .

The ocean floor is like a giant tape recorder

Latitude is deduced from magnetic colatitude θ

$$\tan I = 2 \cot \theta.$$

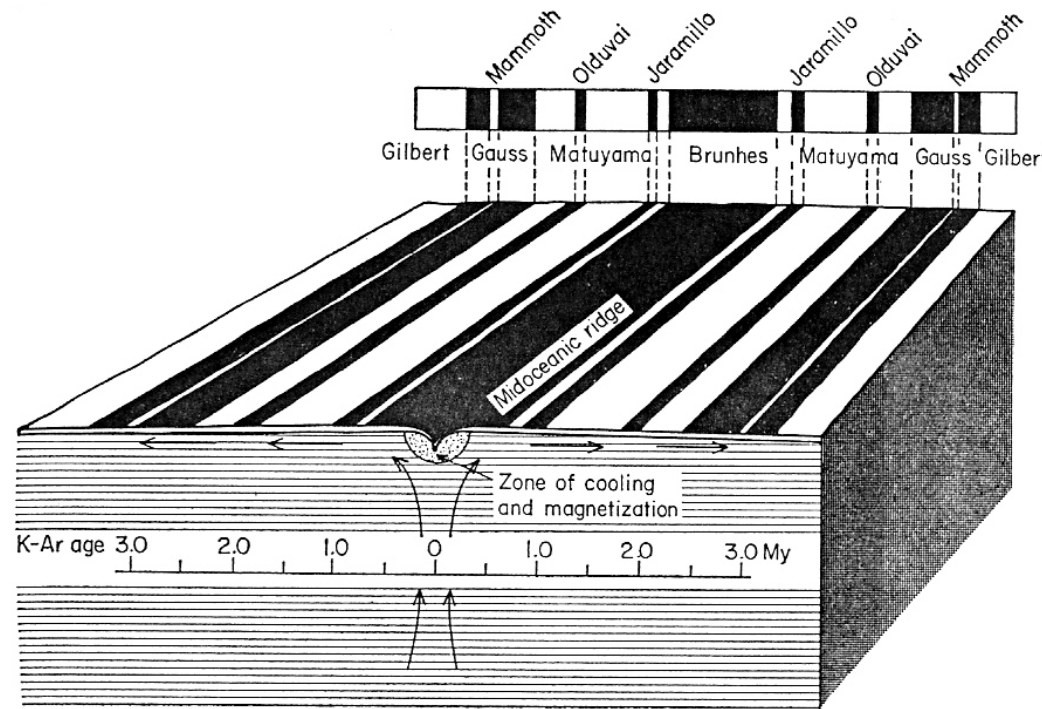
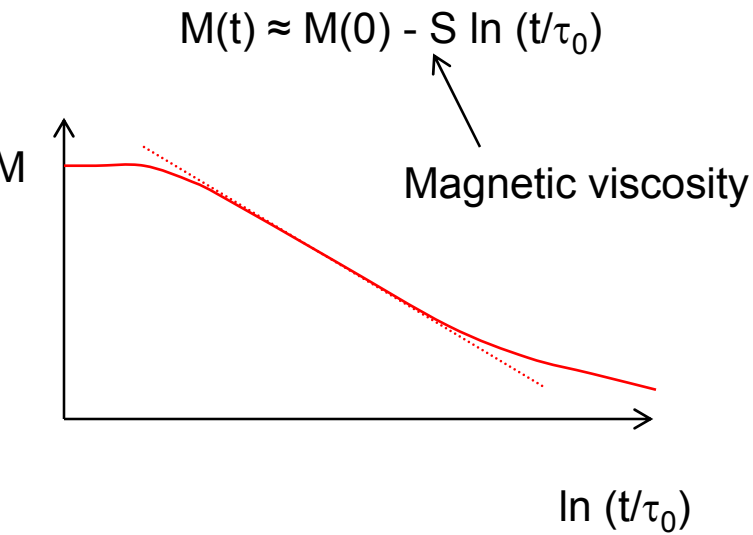


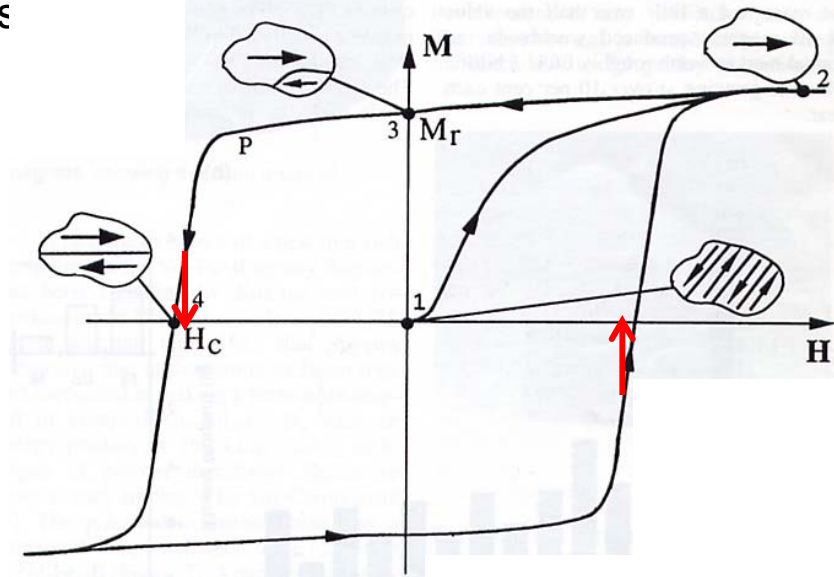
Fig 8.4. Schematic representation of plates separating at a mid-ocean ridge.

3.1 Magnetic viscosity

The stable state of a bulk ferromagnet is a multidomain state with no net magnetization. The magnetic states around the hysteresis
 Metastability is most evident near the coercive field.



$$M(t) = M(0) \int P(\tau) \exp(-t/\tau) d\tau$$



3.2 Ferrofluids

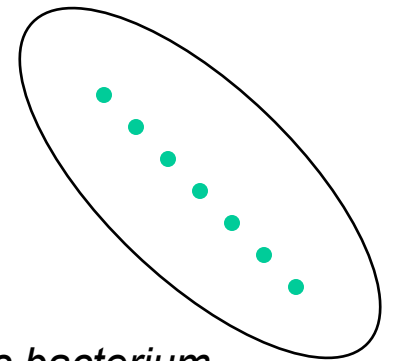
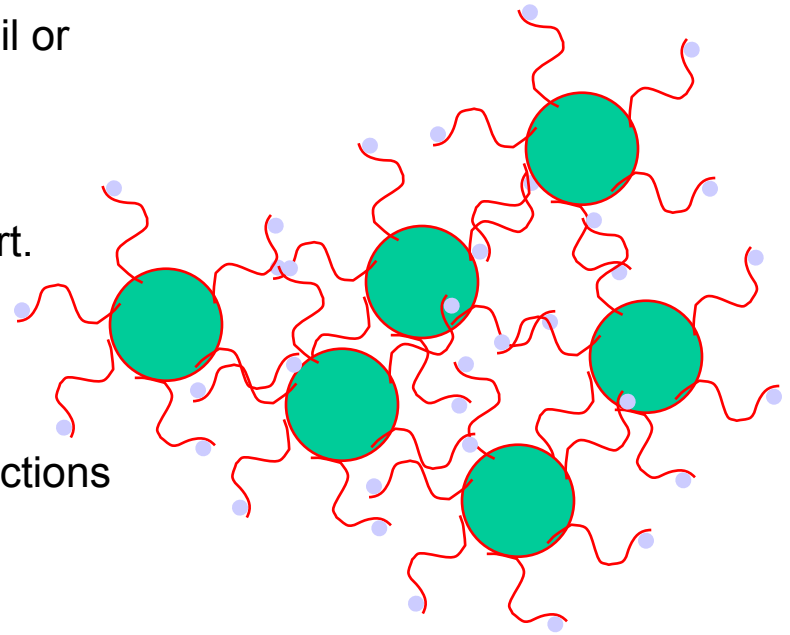
Colloidal suspensions of superparamagnetic particles in oil or water. They behave like ferromagnetic liquids.

Particles are coated by surfactant, which keeps them apart.

A stable ferrofluid must be stable under the influence of gravity, and under the influence of the dipole-dipole interactions (may suppress the superparamagnetic fluctuations) .

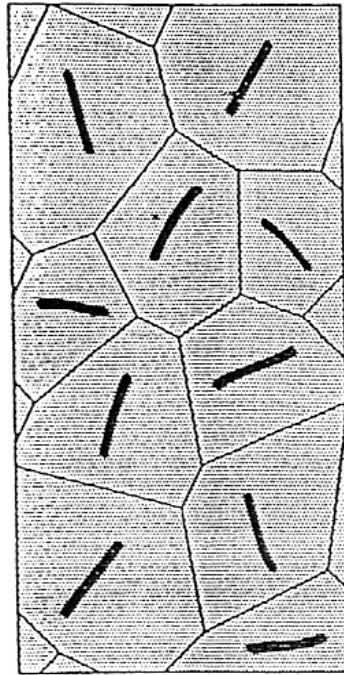
Uses: vacuum bearings, separation by floatation (effective density depends on applied magnetic field; external susceptibility ≈ 3).

Magnetorheological fluids have no surfactant; dipole interactions and viscosity are controlled by an external magnetic field.

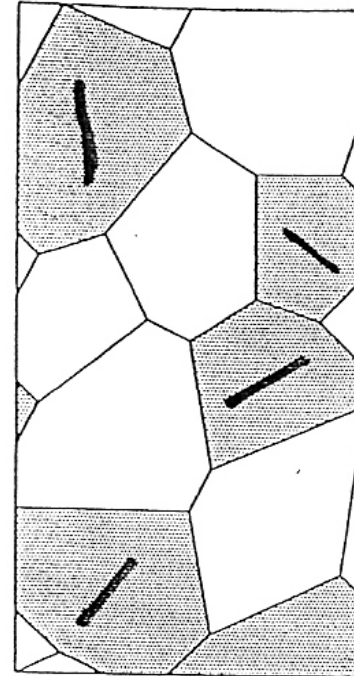


Magnetotactic bacterium

4 Bulk nanostructures



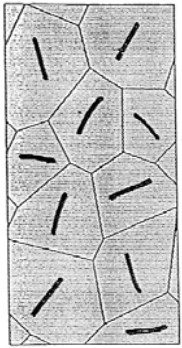
(a)



(b)

Single and two-phase magnetic nanostructures. The easy axis in the harder phase is marked.

The nanostructures may be exchange-coupled across the grain boundaries.



(a)

4.1 Single-phase nanostructures

Exchange-averaging of the anisotropy arises when

- Crystallites are single-domain, with a crystallite size $D \ll$ domain wall width δ ,
- There is exchange coupling across the grain boundaries.

Exchange averaging is effective over the length scale δ_w

A volume δ_w^3 includes $N = (\delta_w/D)^3$ crystallites.

Total anisotropy of the volume obtained by adding

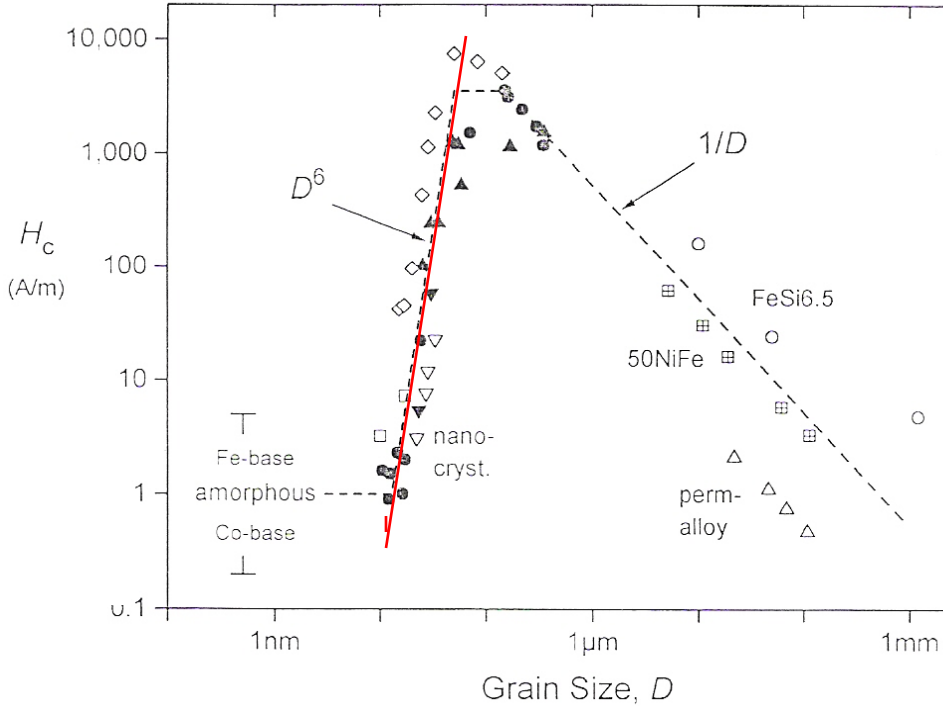
$$\langle K \rangle \approx K_1(D/\delta_w)^{3/2}$$

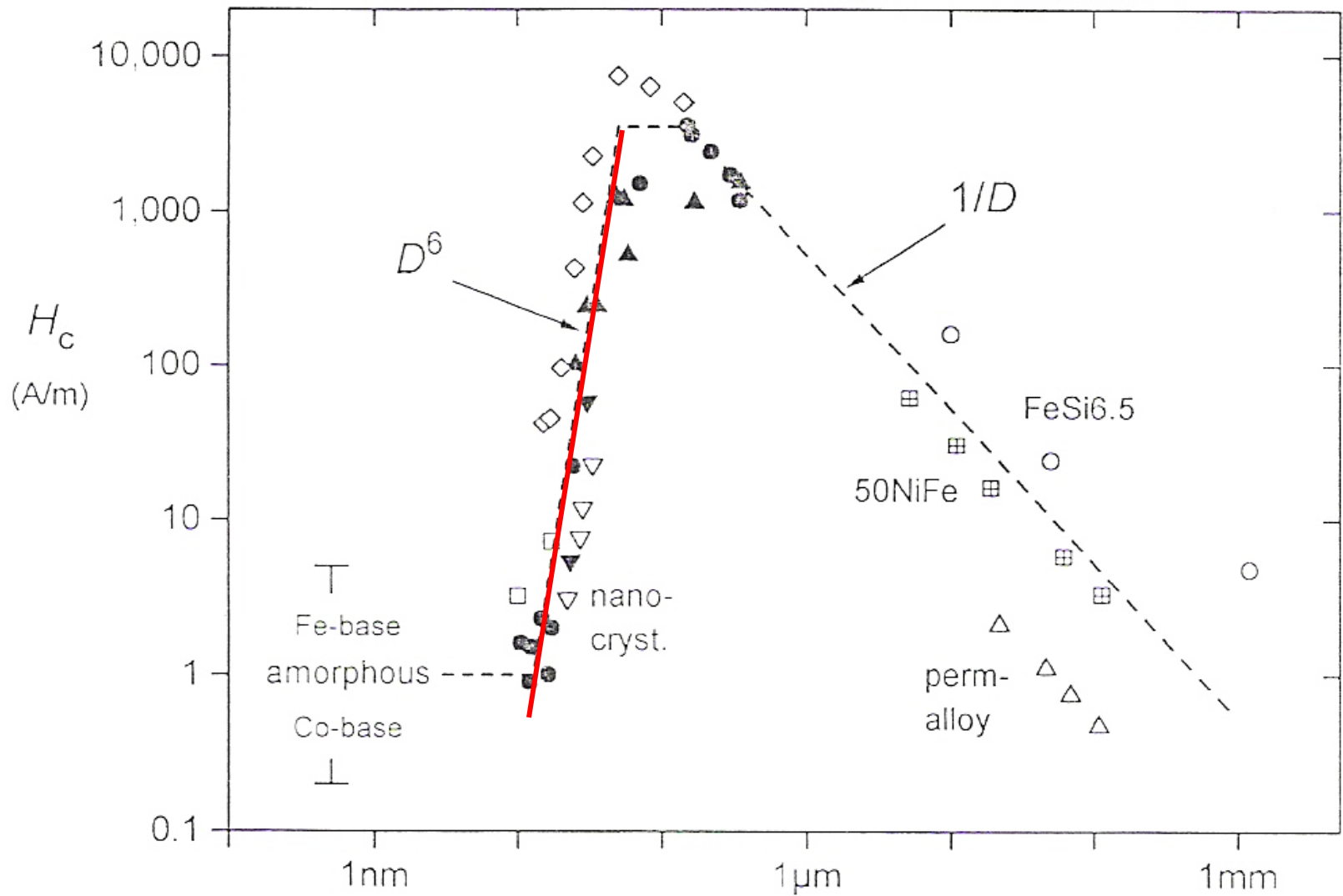
Use this consistently in the expression for the wall

$$\delta_w = \pi\sqrt{(A/K)}; \text{ hence}$$

$$\langle K \rangle = K_1^4 D^6 / \pi^3 A^3$$

Coercivity vs. grain size for a range of soft magnetic materials.





Coercivity vs. grain size for a range of soft magnetic materials.

4.2 Two-phase nanostructures

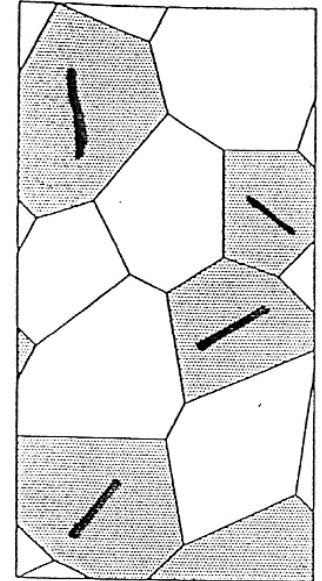
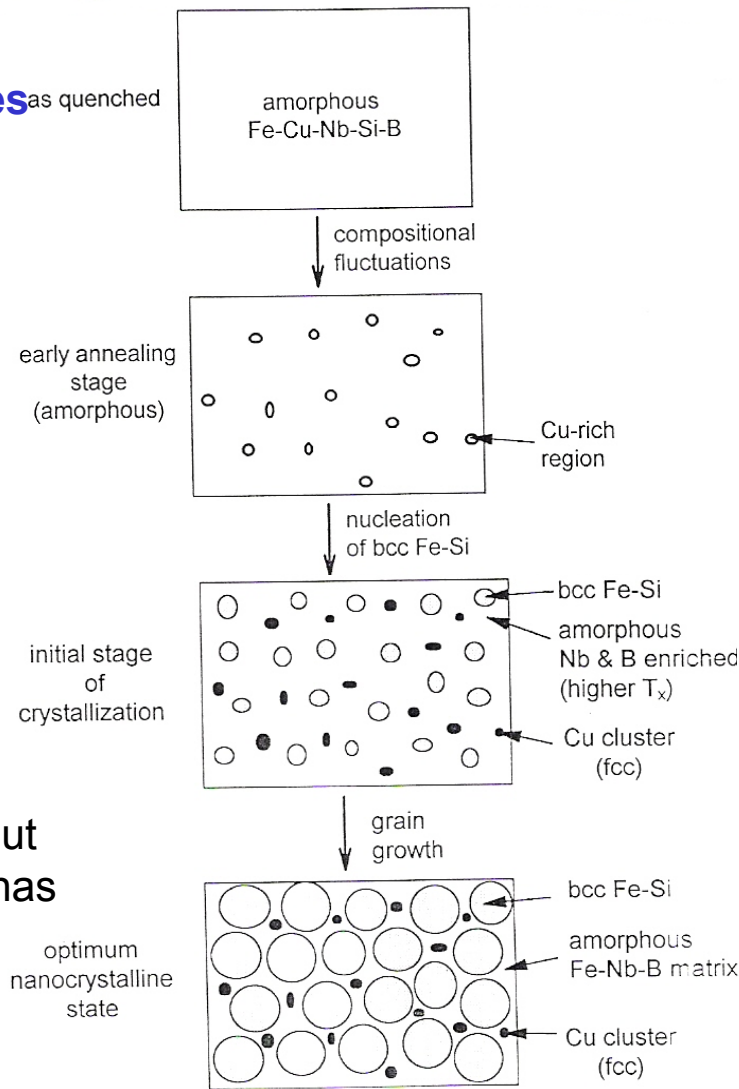
Preparation of *soft/soft* nanostructures

Crystalline particles, in an amorphous matrix

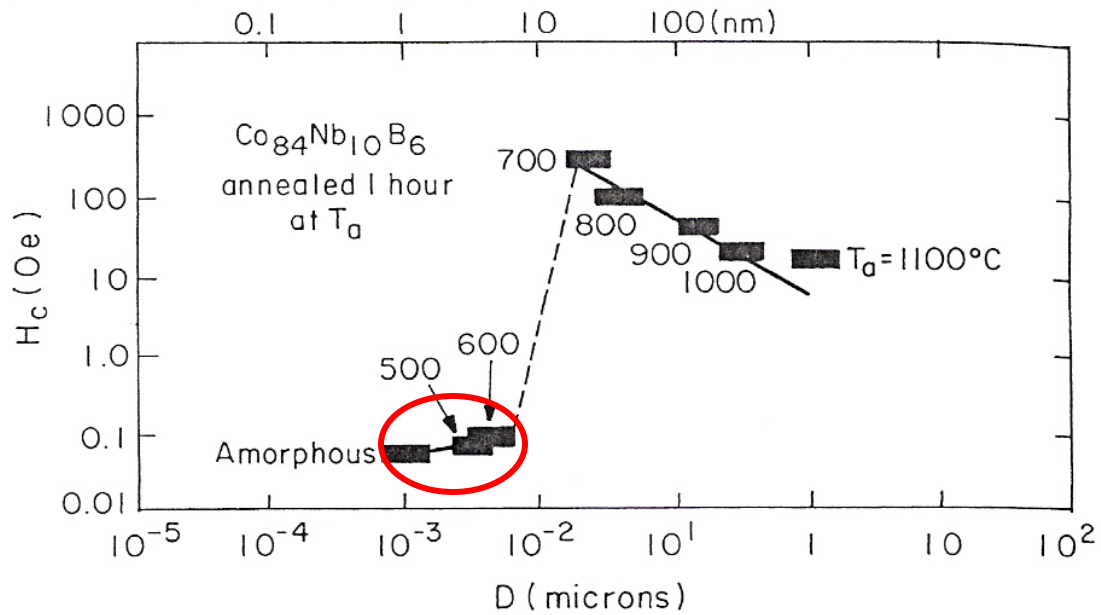
If the volume fraction of the crystalline phase is v , the anisotropy is given by:

$$\langle K \rangle = v^2 K_1^4 D^6 / \pi^3 A^3$$

Method can be used to cancel out the magnetostriction if $\lambda_{\text{crystalline}}$ has opposite sign with respect to $\lambda_{\text{amorphous}}$



crystallization of amorphous Fe-Cu-Nb-Si-B to obtain a two-phase crystalline/amorphous soft nanocomposite 'Finemet'



Coercivity of a partially recrystallized amorphous Co-Nb-B alloy.

5 Needles and wires

Acicular particles with shape anisotropy are used for magnetic recording (tapes and floppy discs).

Shape anisotropy $K_d = [(1-3N)/4]\mu_0 M_s^2$



For a wire $N = 0$. Hence, $K_d = \mu_0 M_s^2/4$. The maximum limit of the coercivity is the anisotropy field $H_k = 2K_d/\mu_0 M_s = M_s/2$

For a true permanent magnet (one that can be made into any desired shape) $H_c > M_s/2$. Shape anisotropy is therefore not enough to make a truly permanent magnet.

Example CrO_2 tapes; particles are $30 \times$

Alnicos are two-phase
CoFe/NiAl nanostructures where
the permanent-magnet
properties are due to the shape
anisotropy of the CoFe acicular
regions.

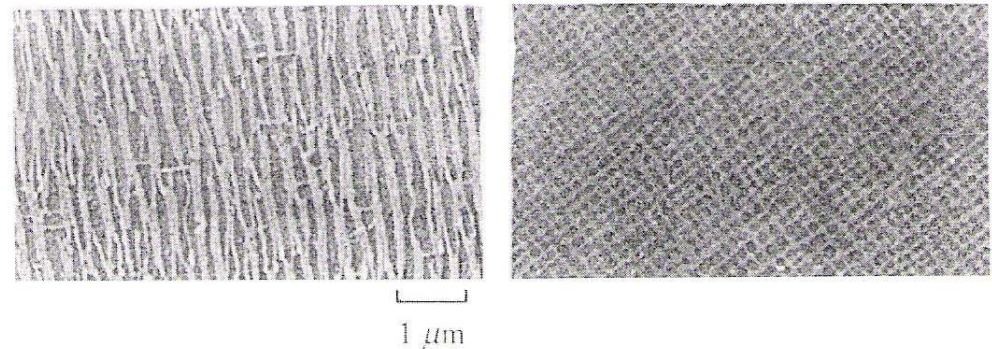


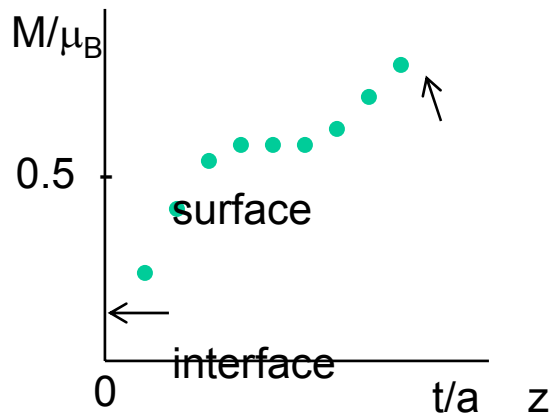
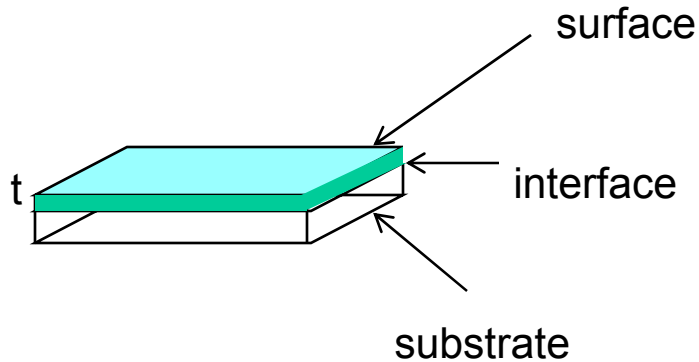
Fig 8.12 The microstructure of an aligned Alnico magnet, showing Fe-Co needles embedded in a nonmagnetic Ni-Al matrix.

6 Thin films

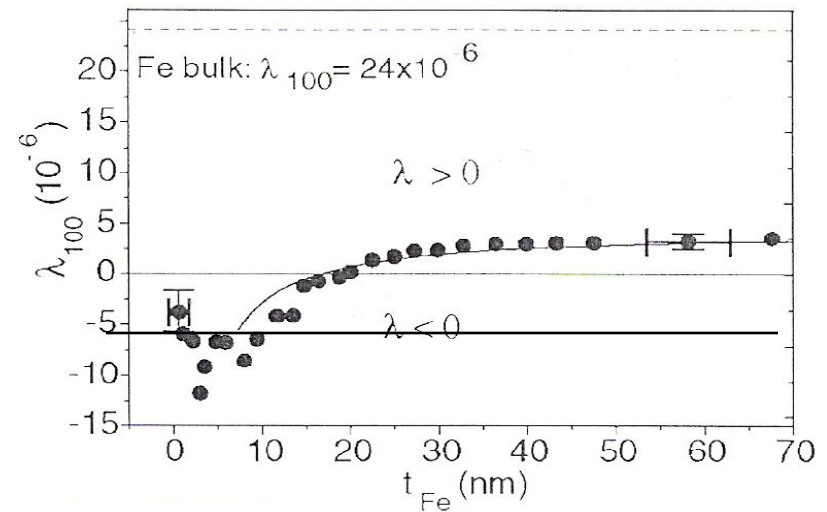
Intrinsic magnetic properties may be different in thin films than in the bulk – Curie temperature, magnetization, anisotropy, magnetostriction.

Reasons are:

- the numbers of surface ions t/a , and interface ions t/a
- difference in lattice parameters in *epitaxial* films; \approx few %



Moment in an 8-monolayer film of Ni on Cu



Magnetostriction of iron thin films

6.1 Magnetization and Curie point

Dramatic changes in magnetic properties may be found in very thin films.

- Band narrowing at the surface may cause ferromagnetism in some d-elements, e.g. V, Pd (Stoner criterion)
- Films grown epitaxially on their substrate can have different lattice parameters or crystal structure to the bulk, hence potentially very different magnetic properties. E.g. Fe grown on Cu (fcc) may be nonmagnetic or antiferromagnetic, according to the substrate temperature.
- Films with different orientations may have different moments. e.g. Surface layer of 100 iron has atomic moment of $3.0 \mu_B$, 110 has $2.6 \mu_B$
- Hybridization with the substrate usually reduces the moment, and may even change sign of J'

Curie point may decrease in thin films due to weakened exchange of surface / interface layers, or it may increase in some cases because of band narrowing.

Note that a *uniformly magnetized thin film produces no stray field*

$$B_{\perp} = 0; \quad H_{\parallel} = 0$$

6.2 Anisotropy and domain structure

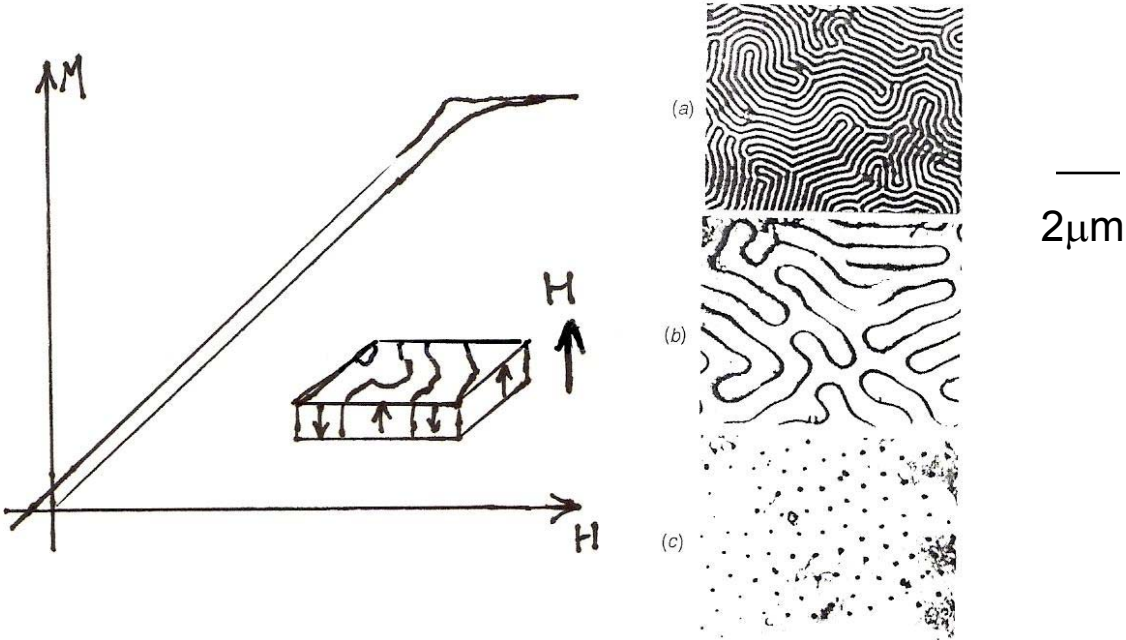


Fig 8.15. Magnetization and domain structure of a thin film with perpendicular anisotropy

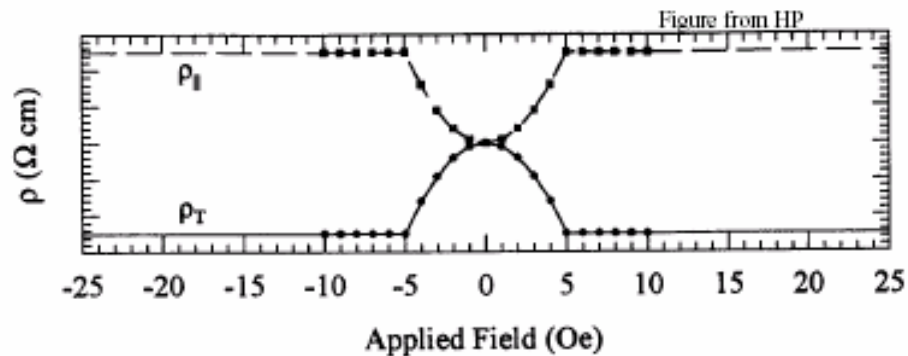
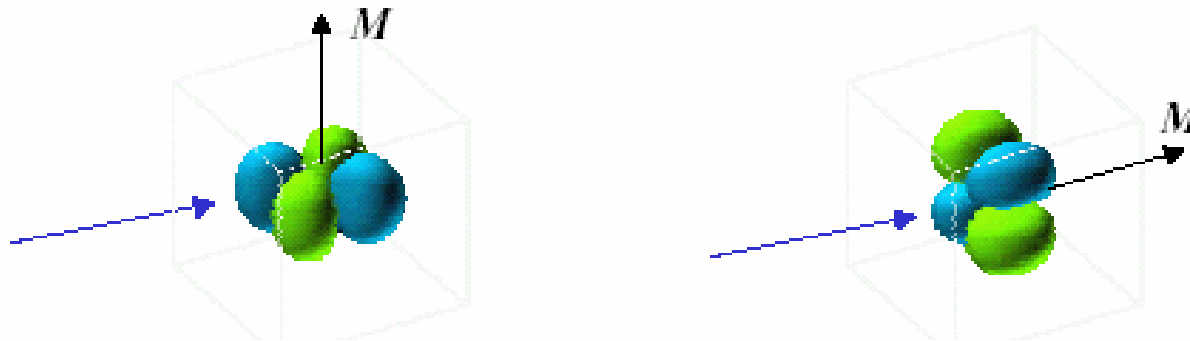
3 Anisotropic magnetoresistance

Feature of magnetic materials

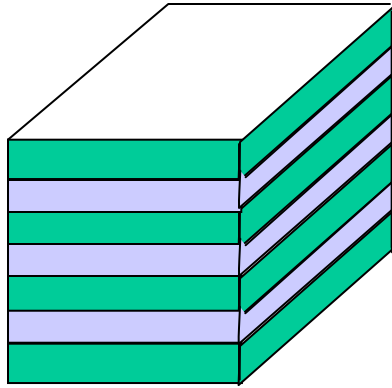
AMR originates from spin-orbit coupling.

It can be positive and negative and its magnitude depends on the scattering cross-section that is presented to conduction electrons by the anisotropic charge distribution of the atoms.

Order of magnitude 1 %.



7 Thin film heterostructures



Practical devices are composed of > 10 layers made of six different metals; Some are very thin e.g. Ru, Co-Fe, AlO_x 1 - 2 nm.

7.1 Giant magnetoresistance

Requires multilayer structures of alternating magnetic and nonmagnetic materials

Origin in spin-dependent electron transmission at interfaces and spin-dependent conductance in ferromagnetic layers

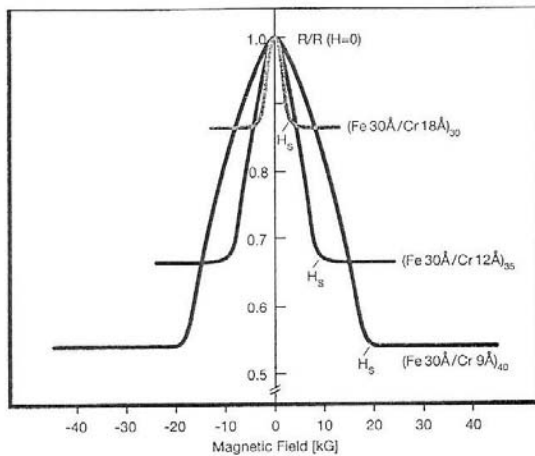
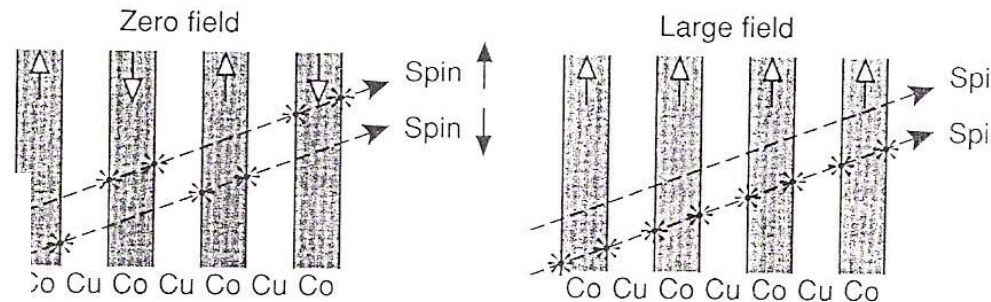
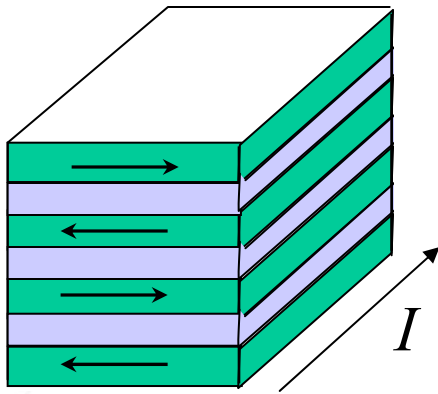


Fig 8.25. Giant magnetoresistance in Fe/Cr multilayers.



Illustration of the derivation of Eq 8.8. 99.

7.2 Indirect exchange coupling

The sign of indirect exchange coupling (RKKY interaction) in multilayers with alternating ferromagnetic and nonmagnetic layers oscillates with the thickness of the nonmagnetic layers.

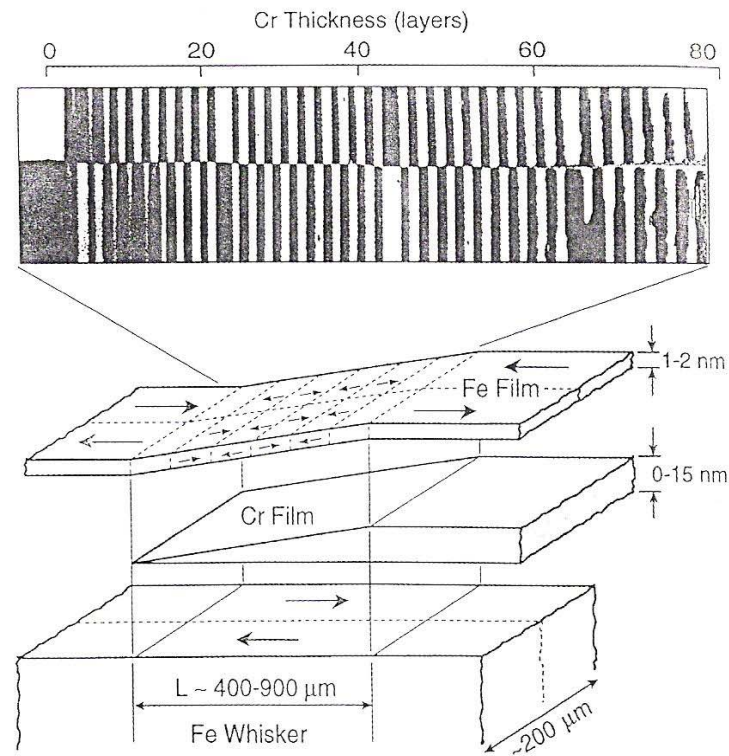


Fig 8.18 An experiment which demonstrates the oscillating spin polarization as a function of spacer thickness.

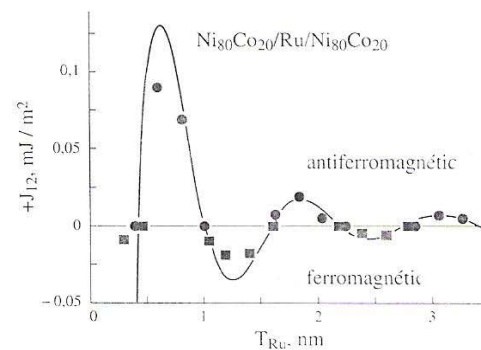


Fig 8.19. Oscillating exchange coupling in Co/Ru/Co trilayer

7.3 Dipolar coupling

There is no dipolar coupling between perfectly-smooth uniformly-magnetized ferromagnetic layers (no stray field).

However, in an actual multilayer the interface are rough and this creates a dipolar coupling field.

If the roughness is correlated the dipolar coupling (orange-peel coupling) is ferromagnetic.

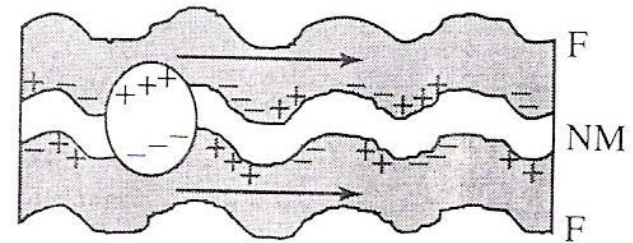


Fig 8.21. The orange-peel effect

7.4 Exchange bias

Needed to *fix* the direction of one of layers.

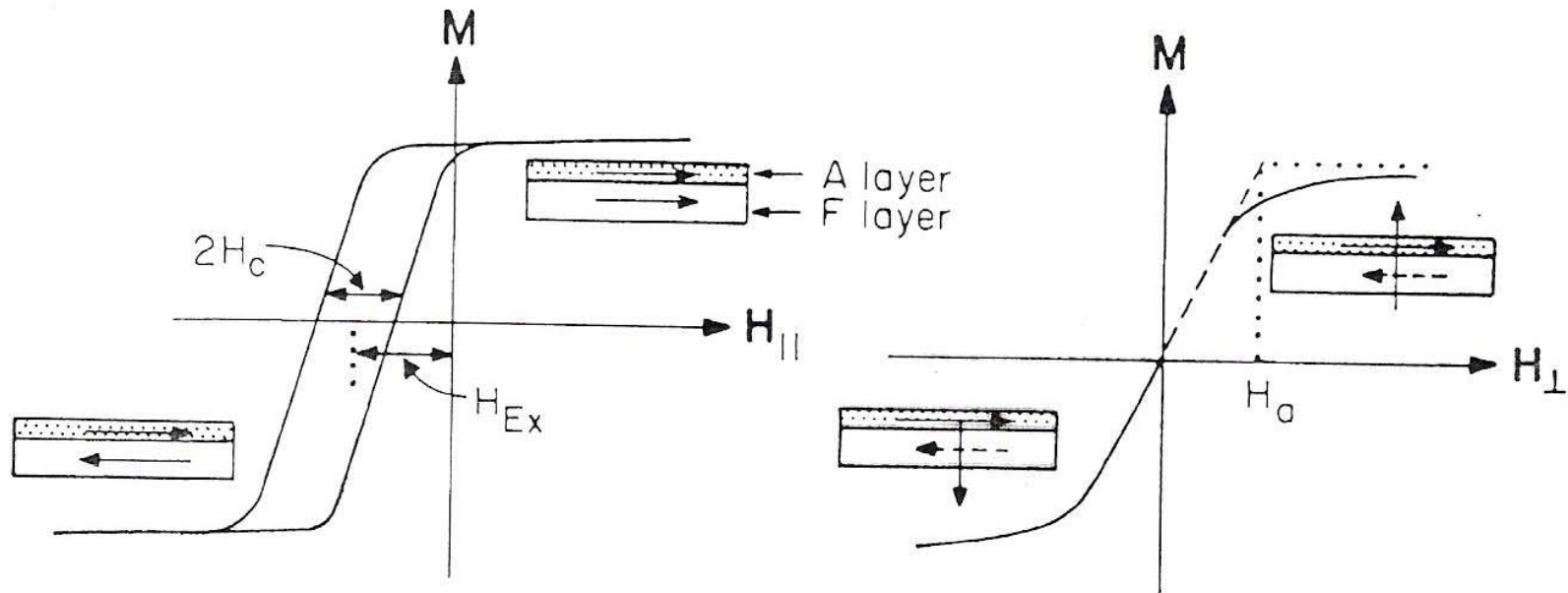


Fig. 8.3 The effect of exchange anisotropy on the hysteresis loop of a ferromagnetic layer coupled to an antiferromagnetic layer. The arrow in the A-layer shows the direction of the exchange field, which is not necessarily the antiferromagnetic axis. The loop on the left is measured with the applied field in this direction; the one on the right with the applied field in the perpendicular direction.

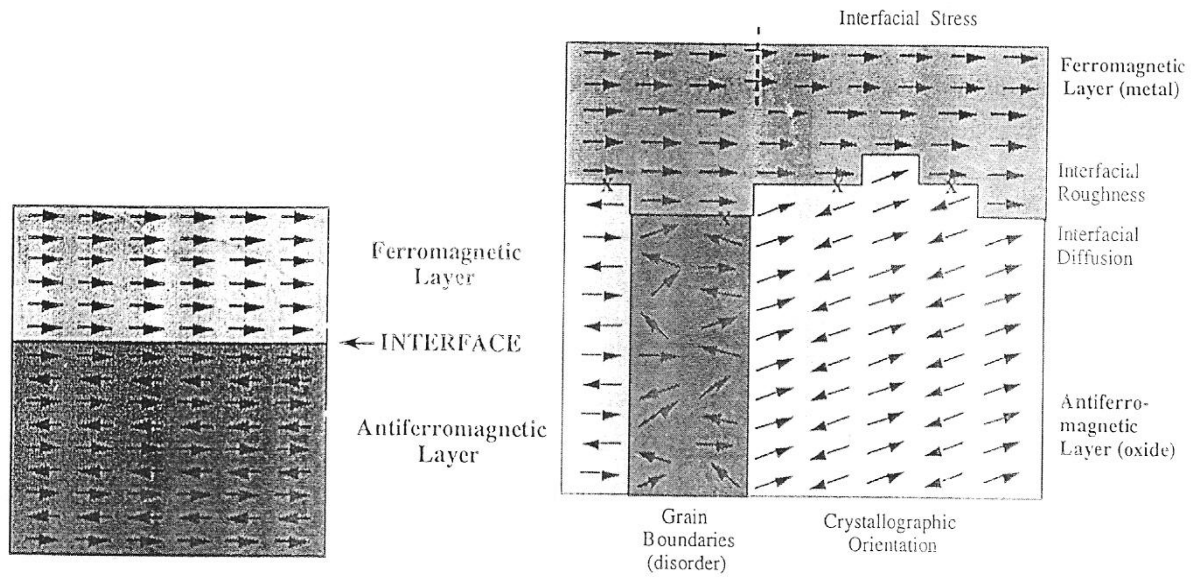


Fig. 8.35 An ideal interface and a real interface.

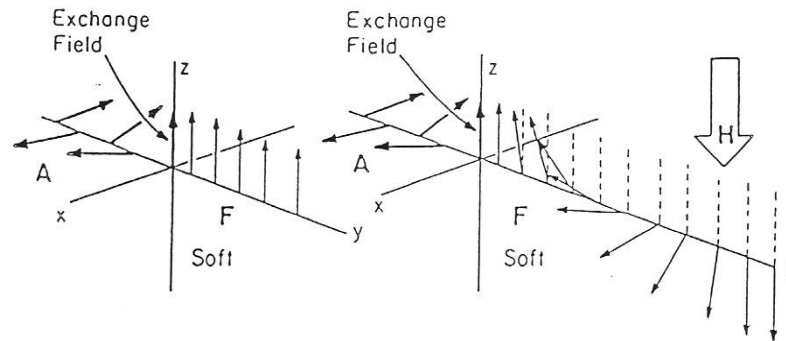
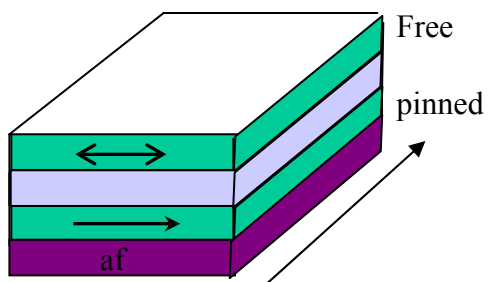
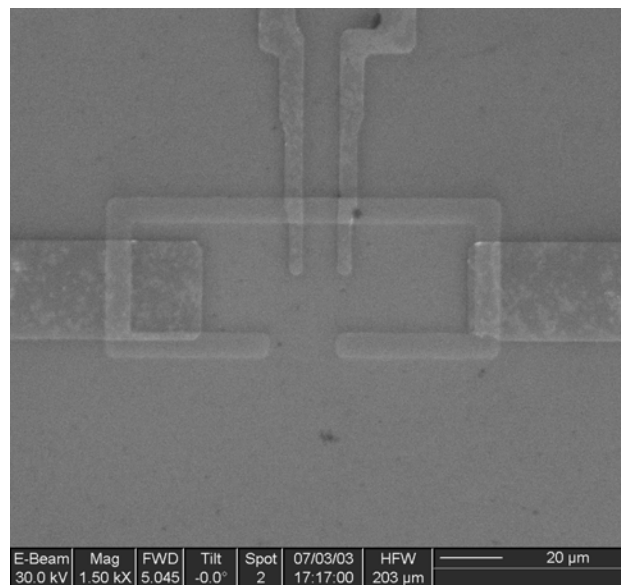
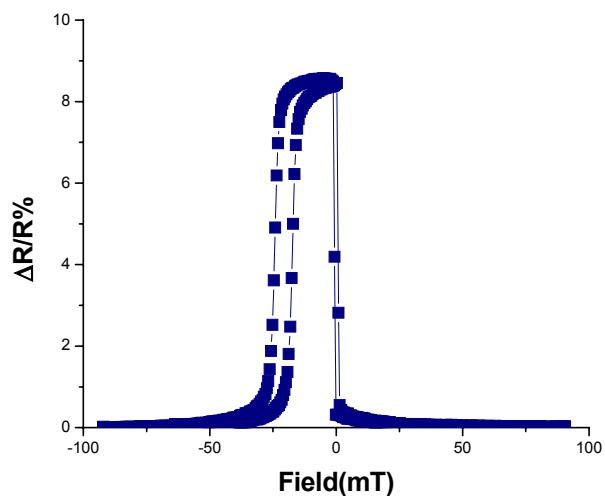


Fig. 8.36. Formation of a domain wall at an interface of a soft ferromagnetic layer exchange-coupled to an antiferromagnet.

7.4 Spin valves



GMR are effect in magnetic spin valves is typically 5 – 20%



Ta(5nm)/NiFe(3.5nm)/CoFe(1.2nm)/Cu(2.9nm)/CoFe(2.5nm)/IrMn(10nm)/Ta(5nm)

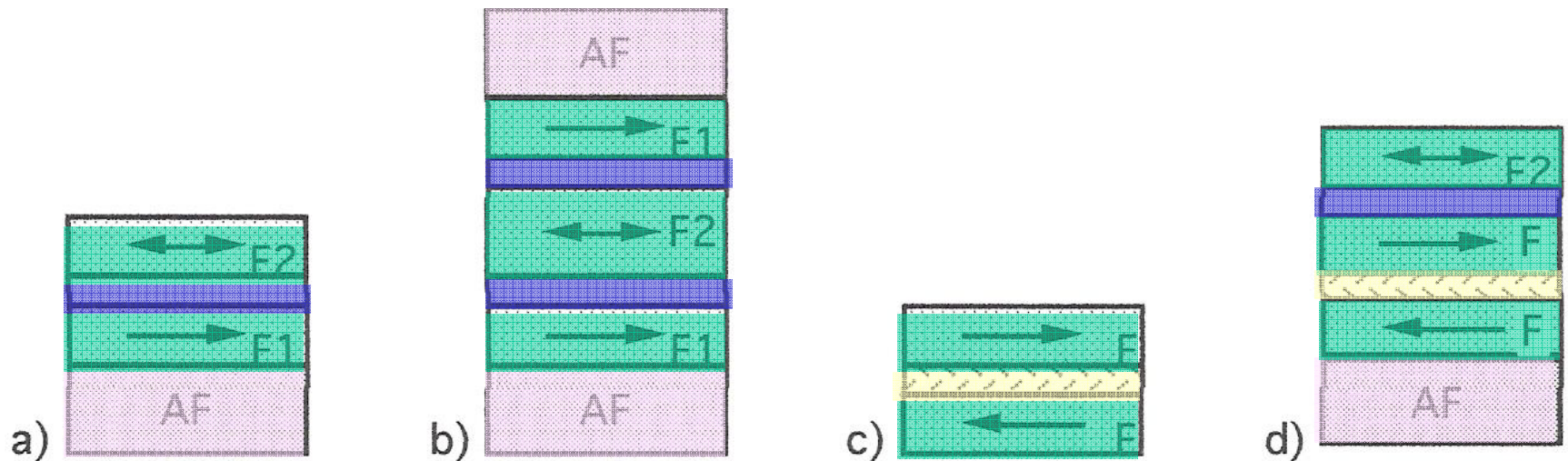
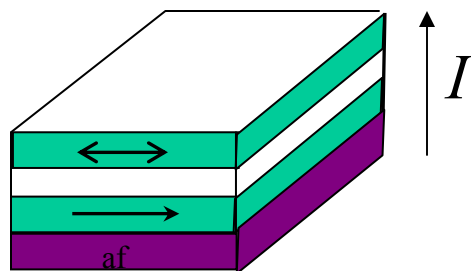


Fig. 8.26. Spin valve structures.: a) simple spin valve with an antiferromagnetic pinning layer, b) double spin valve c) an artificial antiferromagnet, d) a spin valve based on an artificial antiferromagnet. The interfaces between the magnetic layers (F1, F2) and the spacer layer (unshaded) are often decorated with an ultrathin cobalt layer to improve $\Delta\rho/\rho$ for the devices.

7.5 Metal/insulator/metal tunnel junctions



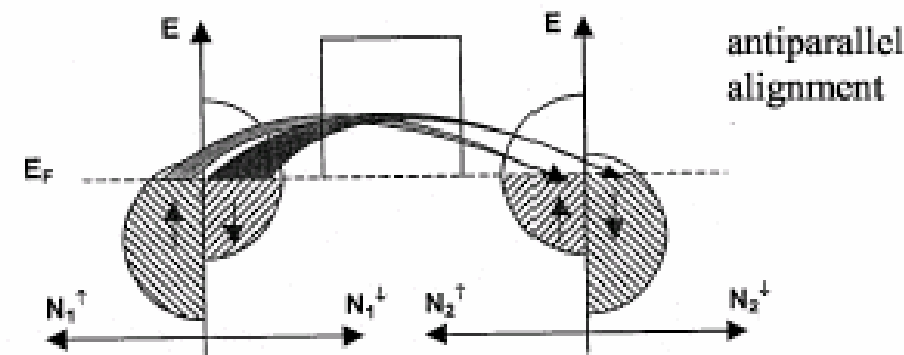
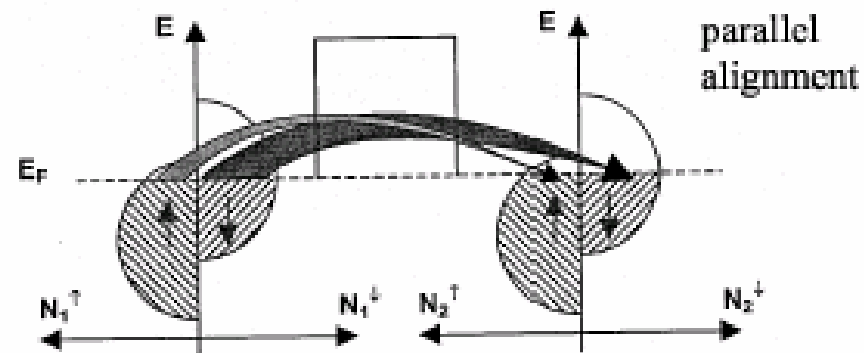
Define spin polarization

$$P = \frac{N_{\uparrow} - N_{\downarrow}}{N_{\uparrow} + N_{\downarrow}}$$

P is about 40% for 3d transition metals

TMR effect typically 50% for AlO_2 tunnel barriers.

Recently (late 2004) MTJs with MgO tunnel barriers exhibit TMR effects up to about 300%.



$$\frac{G_{\uparrow\uparrow} - G_{\uparrow\downarrow}}{G_{\uparrow\downarrow}} = \frac{(1/R_{\uparrow\uparrow}) - (1/R_{\uparrow\downarrow})}{(1/R_{\uparrow\downarrow})} = \frac{\Delta R}{R_{\uparrow\uparrow}} = \frac{2P}{1-P}$$

$$\frac{\Delta R}{R_{\uparrow\downarrow}} = \frac{2P^2}{1+P^2}$$

7.6 Magnetic single-electron devices

Ion-beam or e-beam lithography

- Quantum dots
- Magnetic Coulomb blockade
- Kondo effect

8 Applications

8.1 Magnetic recording

Analog wire recording was invented in 1898 by Valdemar Poulsen.

Analog tape recording using iron oxide particle tape was developed in Germany in the mid-1930s

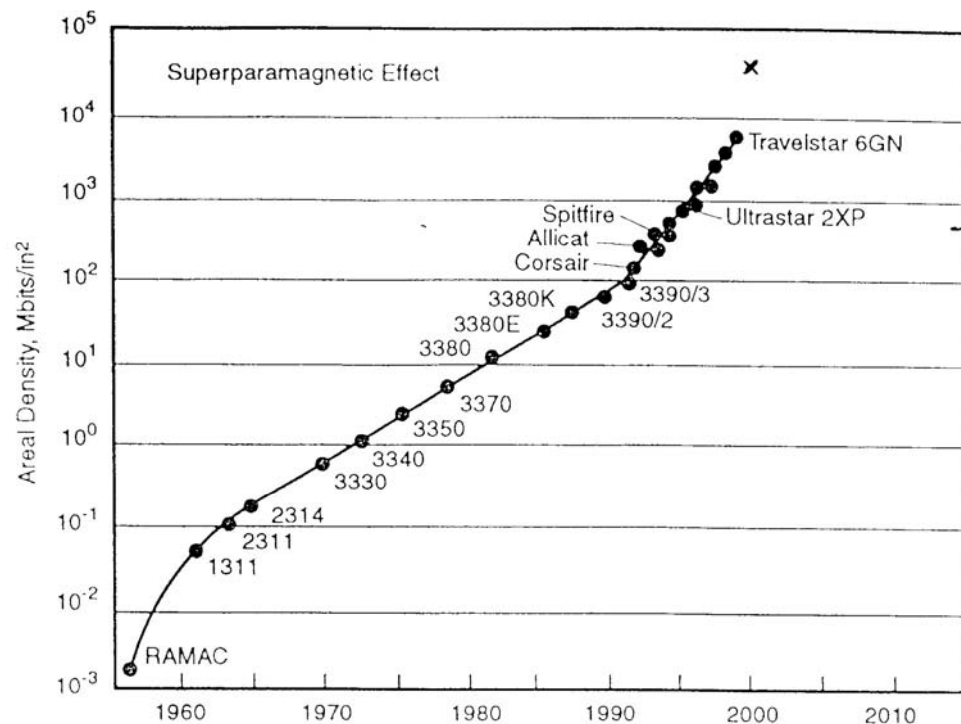
Digital disc recording was introduced by IBM in the mid 1950s

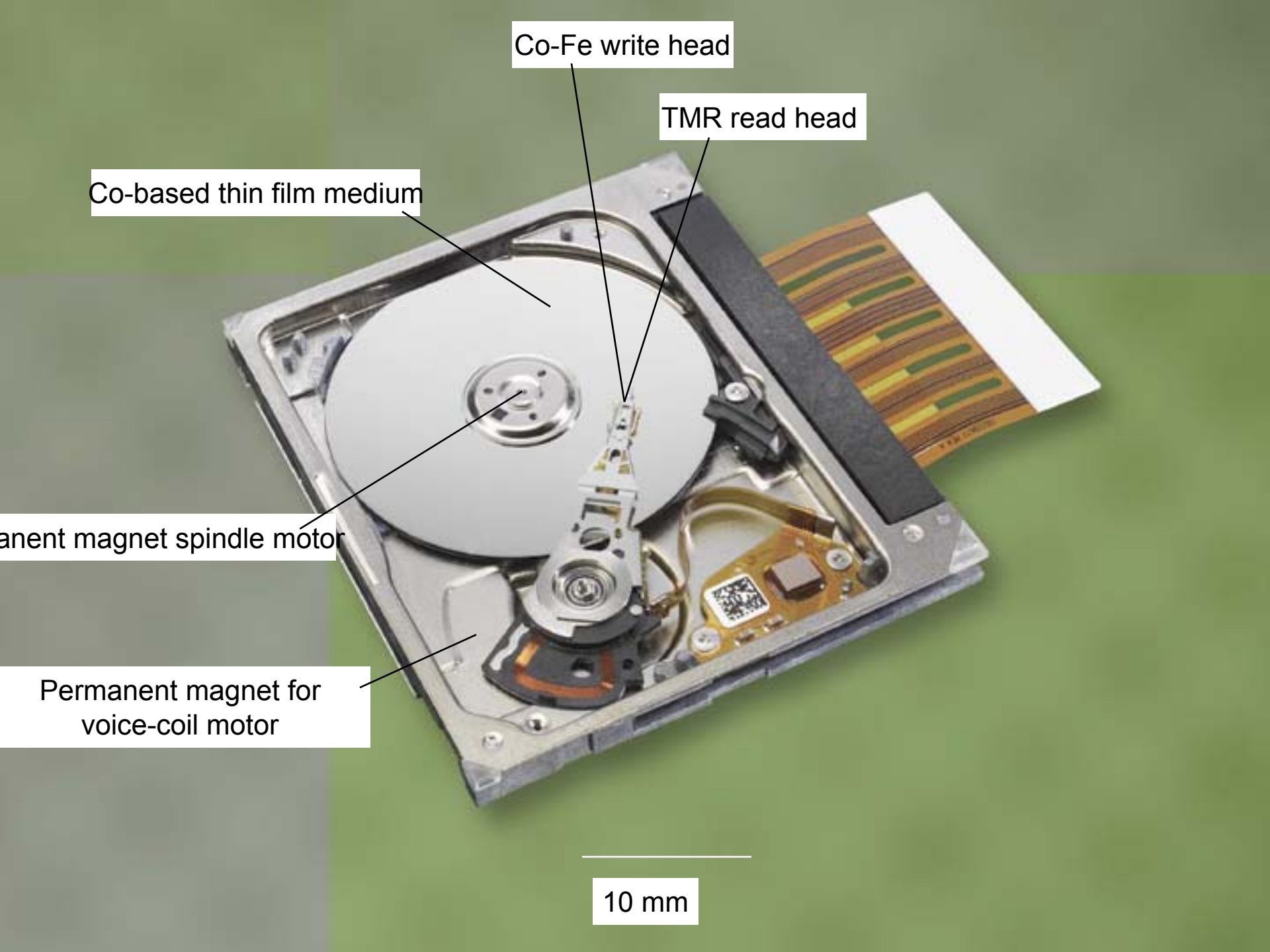
It has been relentlessly perfected over 50 years. Densities have increased by a factor 10^8 to 100 Gb/s per inch. (155 bits μm^{-2}) Data rates are $\sim 1\text{Ghz}$, Fly heights of the head over the disc surface are $\sim 10\text{ nm}$.

Digital and analog recording is a €20 B business, consuming large quantities of ferrite and other semi-hard magnetic materials for recording media, and using sophisticated miniature magnetic circuits in the read and write heads.

The magnetic record is generally in the plane of the medium. Only magneto-optic recording uses perpendicularly-magnetized media at present.

The data are recorded on tracks on the media whose width is determined by the width of the write head.





Co-Fe write head

TMR read head

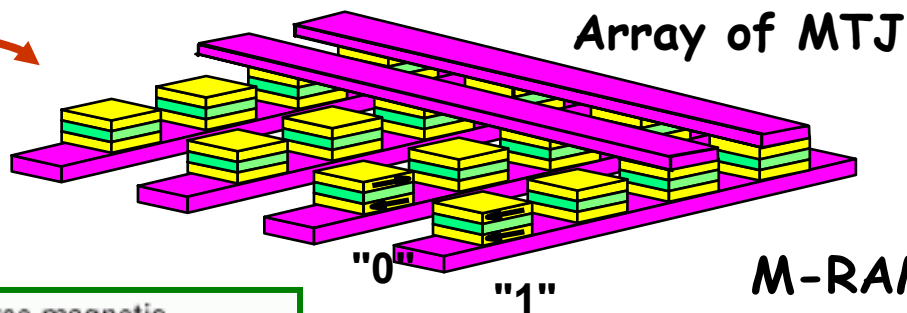
Co-based thin film medium

Permanent magnet spindle motor

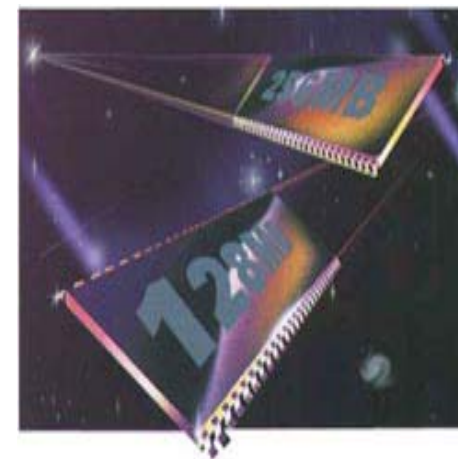
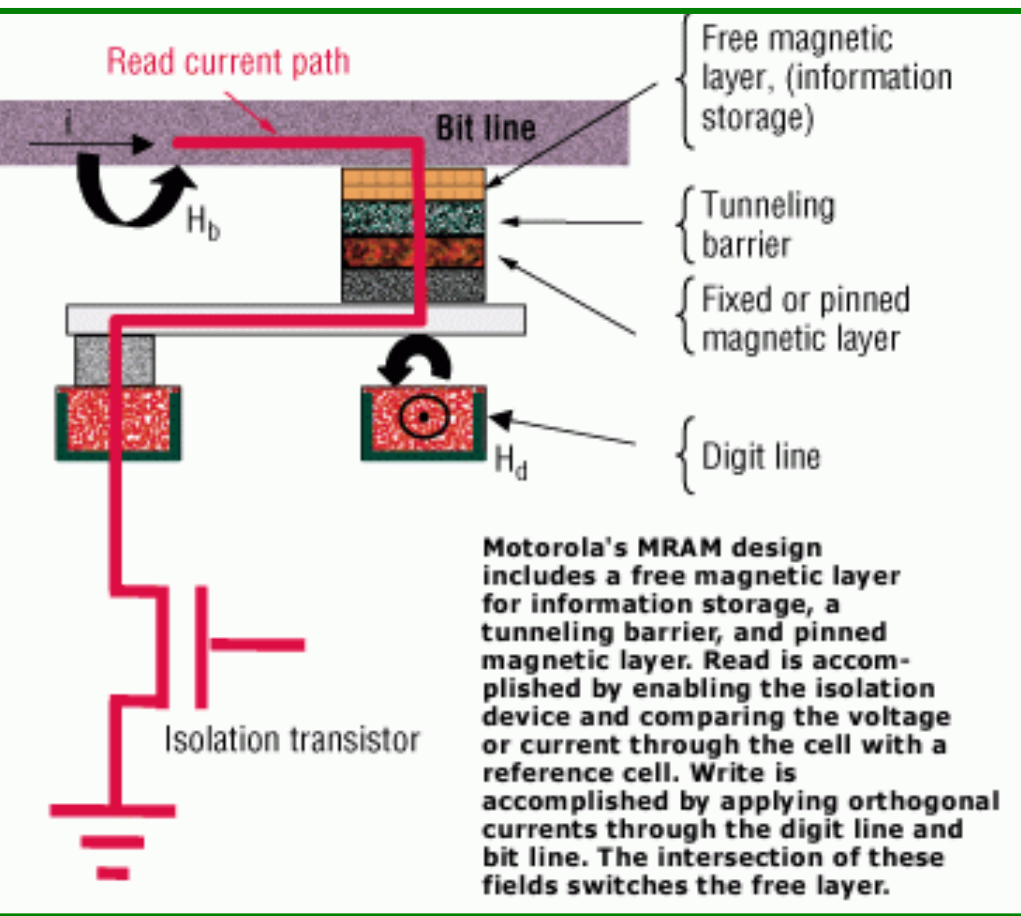
Permanent magnet for voice-coil motor

10 mm

8.2 Magnetic Random - Access Memory (MRAM)



M-RAM



Notebook, PC, etc



Palm Top



Mobile Phone

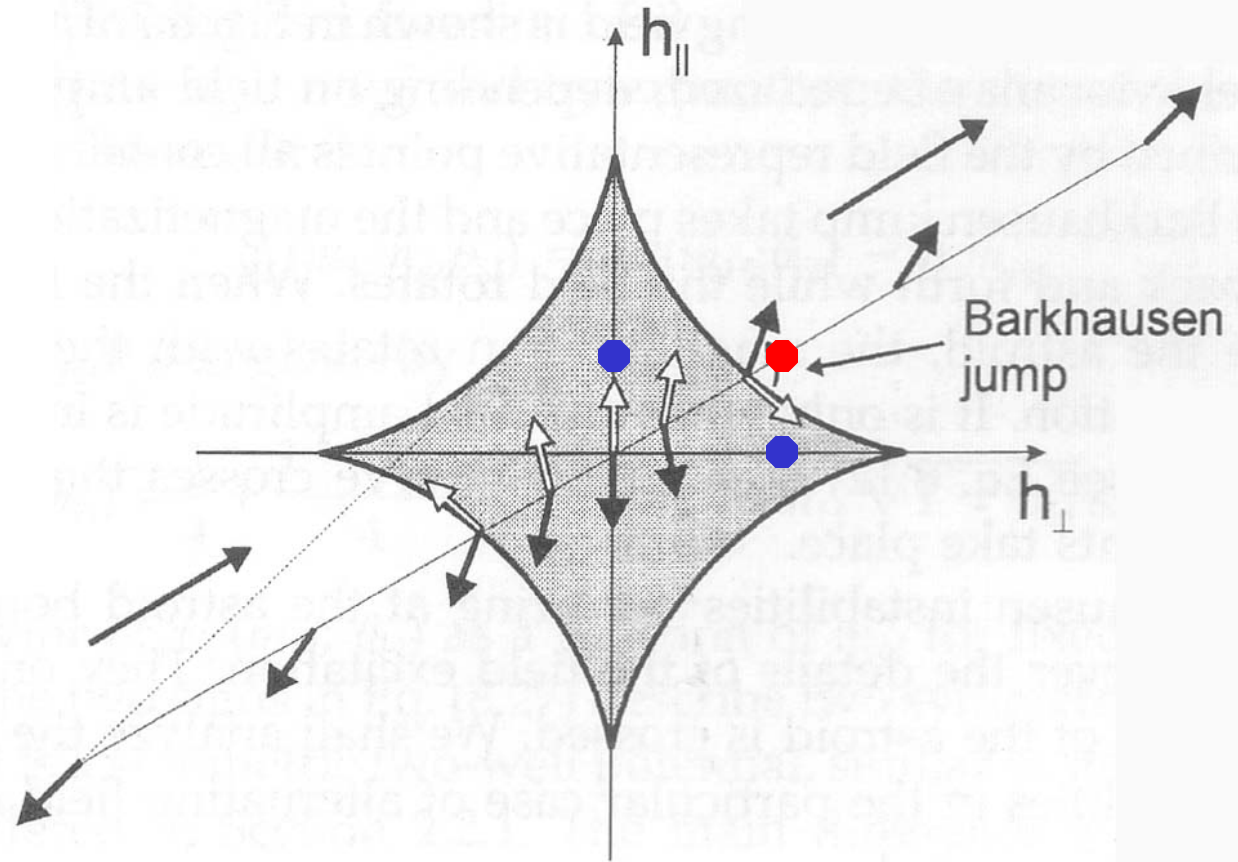
QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Thermally-assisted switching.

Current-induced switching

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.



Half-select concept.

A pulse in the bit line or the word line is not enough to switch the element. *Both* must be present simultaneously.

8.3 Spin Electronics

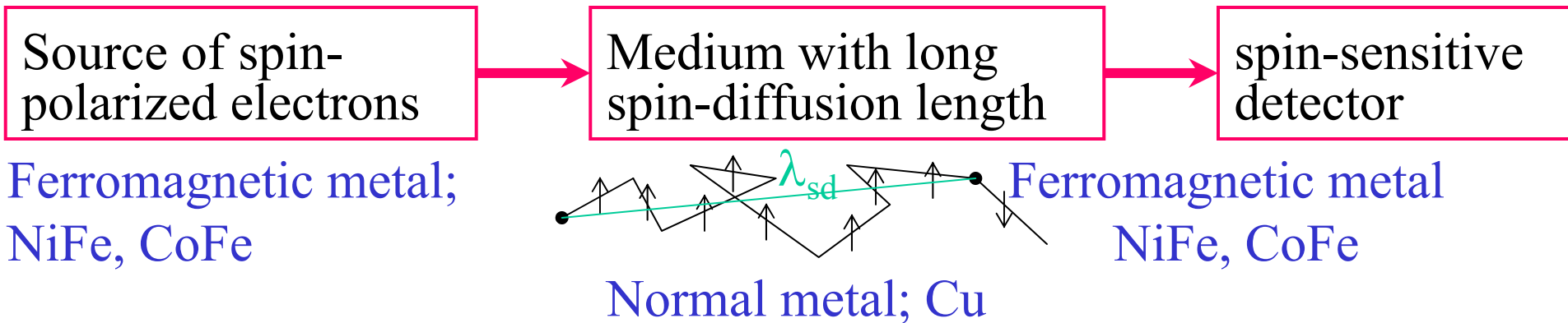
Conventional electronics has ignored the spin in the electron:





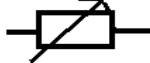


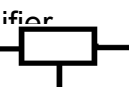
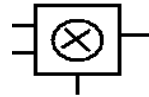
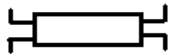
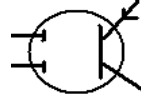
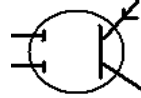



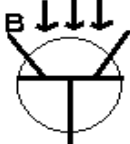
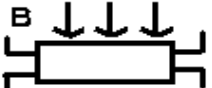
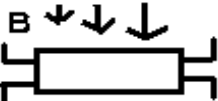
- Code information into the \uparrow and \downarrow channels
- Manipulate the \uparrow and \downarrow electrons independently
- Exploit magnetic and electric fields

$$\uparrow \quad e = 1.6 \cdot 10^{-19} \text{ C}$$

$$m = e\hbar/2m = 1\mu$$

TWO-TERMINAL DEVICES; Magnetoresistors

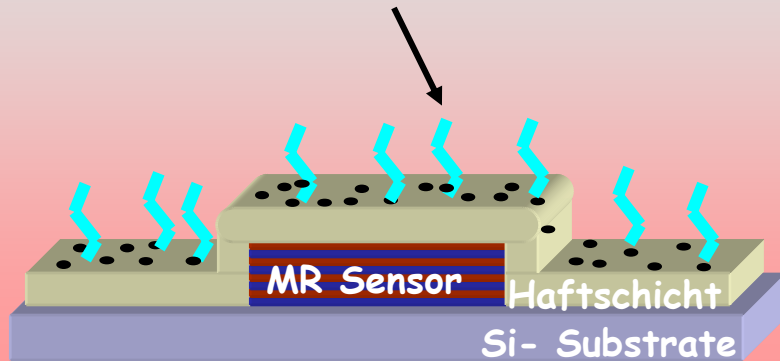


Number of Terminals	2	2+	3 / 3+	4 / 4+
Classical Devices	Switch  Resistor  Diode 	Photodiode  Varistor 	Transistor  Filter  Amplifier 	Wheatstone Bridge  2-gate MOSFET  Tetrode  Multiplier 
Spin Electronic Devices	Spin Switch 	Magnetic switch (MTJ)  Magnetic Photodiode 	Spin transistors 	Hall Probe  Magnetic Gradiometer (bridge) 

8.4 Biochips

1) Immobilisation of target molecules

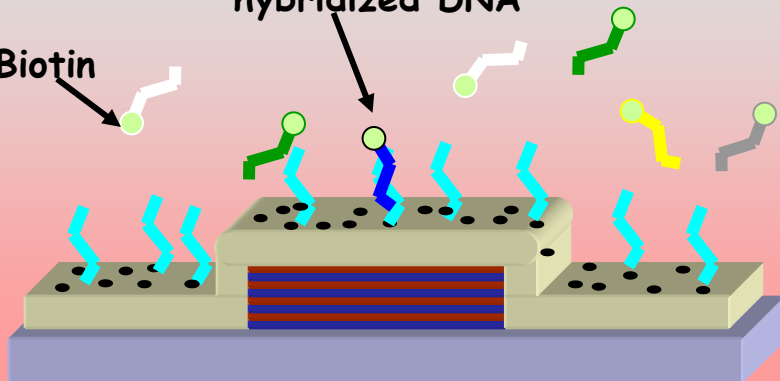
Fixed DNA single strand



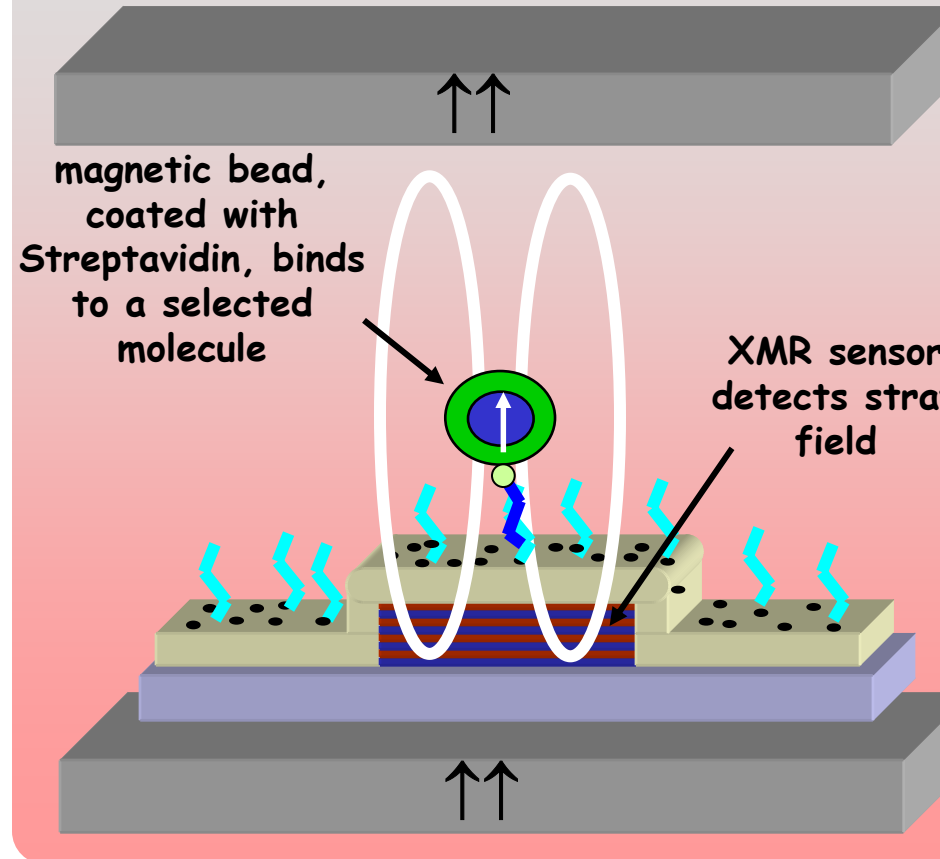
2) Hybridisation of the probe molecules

hybridized DNA

Biotin



3) Hybridisation with beads and detection with MR sensor



Magnetism Conferences

ICM International Conference on Magnetism (IUPAP Support)

- Every three years (IUPAP Support) Recife 2000, Rome 2003, Kyoto 2006, Karlsruhe 2009

APS meeting on Magnetism and Magnetic Materials MMM - Annual

INTERMAG Conference - Annual

Gordon Conference on Nanoscale Magnetism etc.

Suggestion:

IUPAP Meetings: 'Frontiers of Nanoscience' Every two years.

Cross-disciplinary; each supported by two or three commissions.

Nanoscience: Small size 1 - 100 nm;

Complexity

Correlations.

Nanoscience: Science of condensed matter on a scale of 1 - 100 nm where interatomic or interelectronic interactions which impart a complexity which is absent in the bulk.