Capture of Slow Antiprotons by Atoms, Molecules, and Ions

James S. Cohen

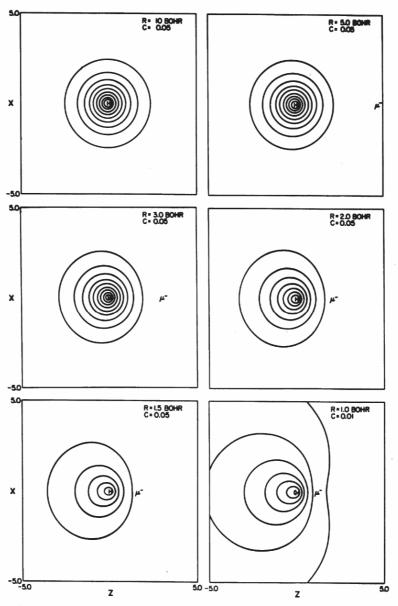
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Workshop on Physics with Ultra-Slow Antiproton Beams RIKEN, March 14–16, 2005

What do we think we know and what can be verified using the new capability with slow antiproton beams?

Recent review (mainly theoretical): J.S. Cohen, Rep. Prog. Phys. 67, 1769-1819 (2004).

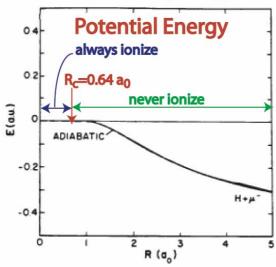
Adiabatic μ^- (or \overline{p}) + H

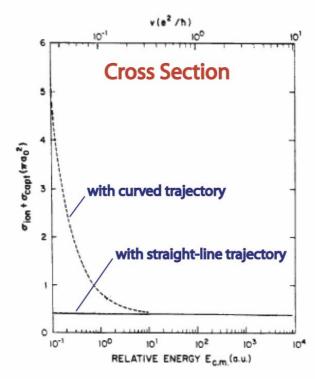


"Adiabatic ionization" occurs at R < $R_c \approx 0.64 a_o$ [Wightman, 1950; Fermi&Teller, 1947]

ADIABATIC IONIZATION

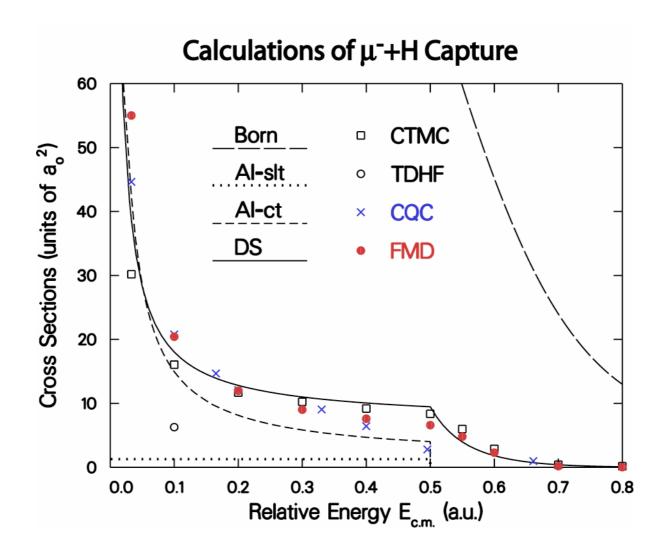
$$x^-+H (x^- = \mu^-, \pi^-, \overline{p})$$

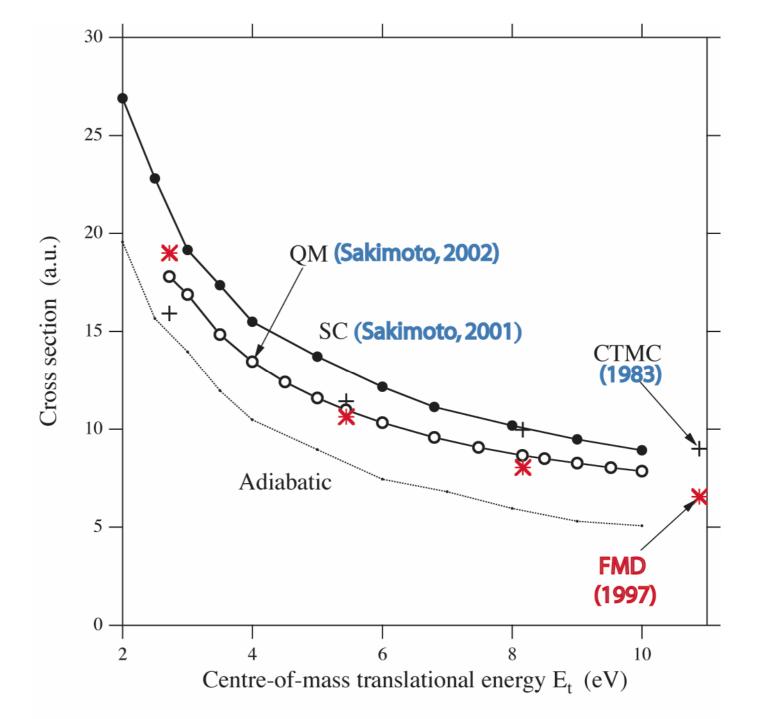


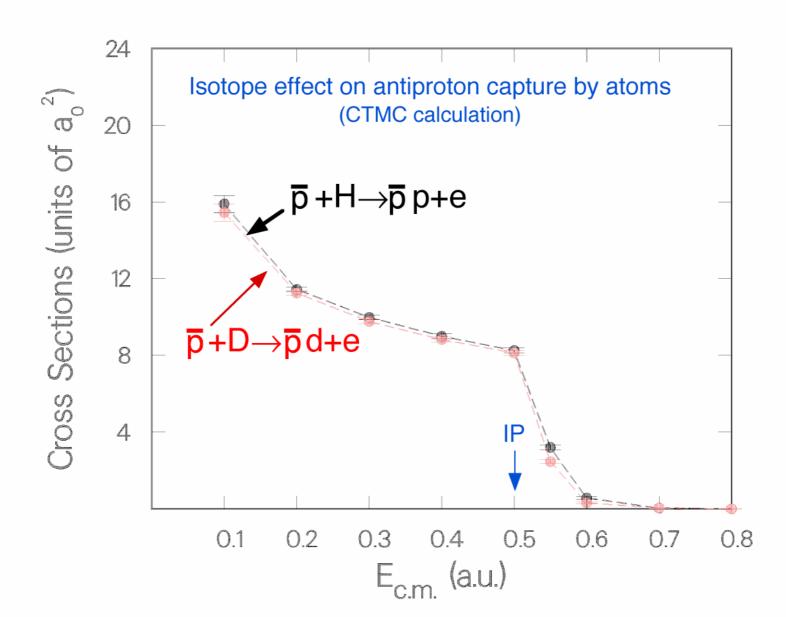


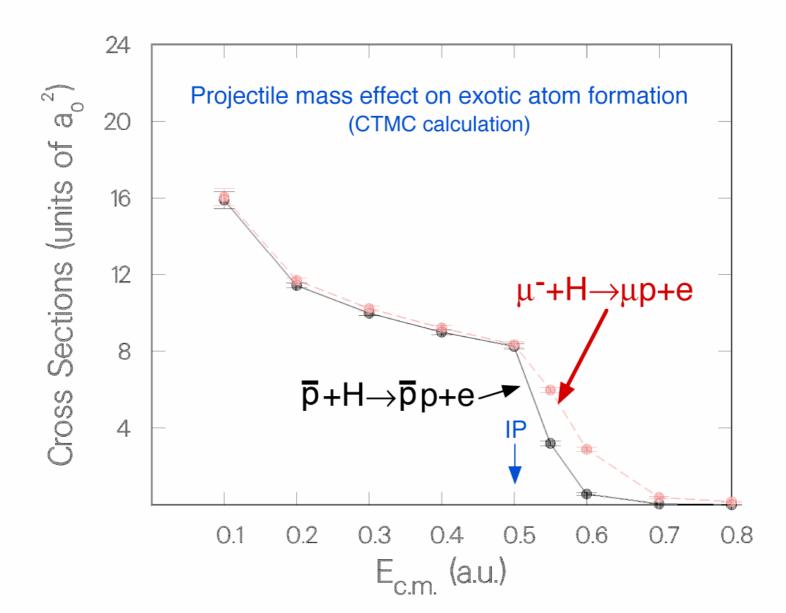
Theoretical methods for X^- capture by the H atom

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1. Adiabatic ionization (AI)
   [Wightman(1950)]
 2. Born / Coulomb-Born (BA, DWA)
   [de Borde(1954) ... Korenman & Rogovaya(1980)]
 3. Diabatic states (DS)
   [Cohen, Martin & Wadt(1981)]
 4. Classical-trajectory Monte Carlo (CTMC)
   [Cohen(1983)]
 5. Time-dependent Hartree-Fock (TDHF) with \mu wave packet
   [Garcia, Kwong & Cohen(1987)]
 6. Classical-quantal coupling (CQC)
   [Kwong, Garcia & Cohen(1989)]
 7. Close coupling (CC)
   [Boukour(1996)]
 8. Fermion molecular dynamics (FMD)
   [Cohen(1997, 1998)]
9. Perturbed stationary states (PSS)
   [Ohtsuki(1999)]
10. Time-dependent semiclassical(trajectory) Schrödinger eq. (SC)
   [Sakimoto(2001)]
11. Time-dependent wavepacket Schrödinger equation (QM)
   [Sakimoto(2002)]
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The Quasiclassical Method

The Fermion Molecular Dynamics (FMD) method is an extension of the Classical-Trajectory Monte Carlo (CTMC) method to multielectron atoms and molecules, obtained by introducing pseudopotentials to represent quantum-mechanical effects.

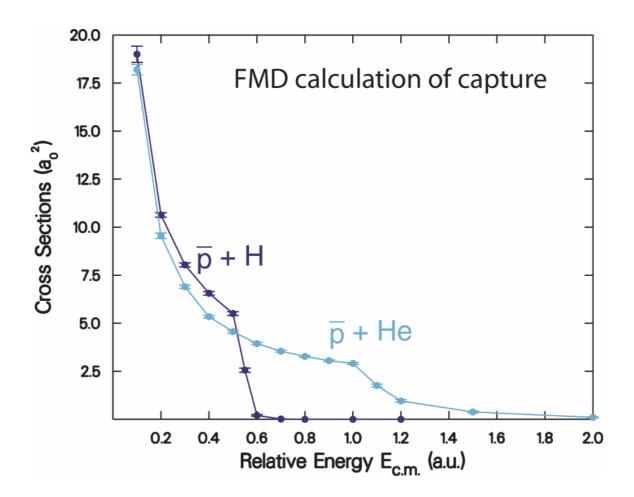
Advantages for capture of heavy $(m \gg m_e)$ negative particles:

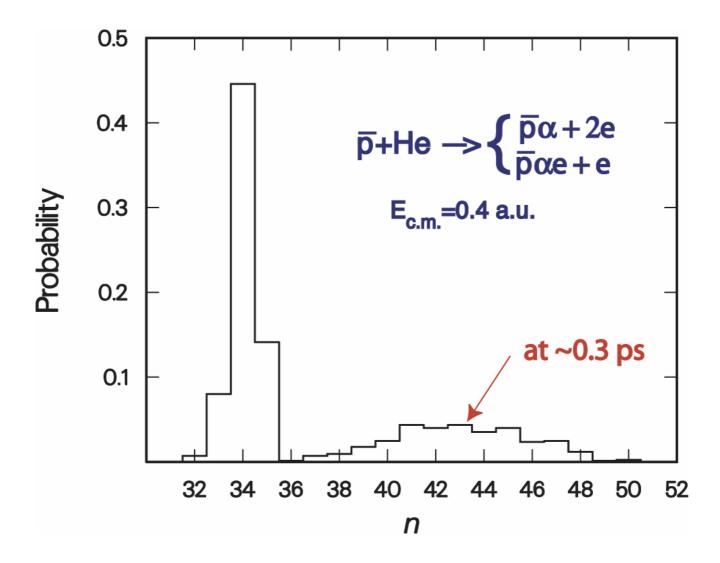
- The classical (Hamilton) equations of motion are easier to solve than the Schrödinger equation.
- Full dynamics is done for all particles. Thus correlation is fully treated and all rearrangement channels (including breakup) are treated on equal footing.
- The electronic continuum occurs naturally.
- ullet The large number of intermediate and final states presents no difficulty. Capture is into high n,l orbitals, so the correspondence principle applies.

Possible weaknesses:

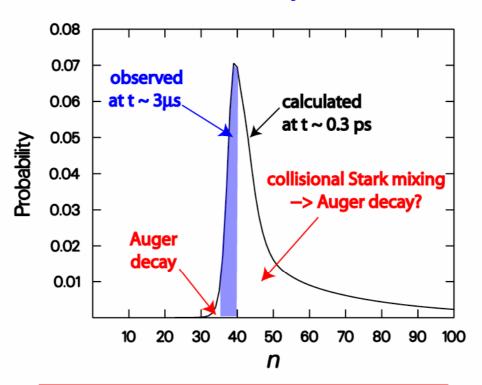
- Correlation may be too strong.
- May not be adequate when the molecular vibrational—rotational spacings are large (as in H₂).
- Not valid for light reactants like electrons and positrons.

Where corroboration is available, the FMD method seems to be rather accurate, but, at the very least, it is useful to obtain an indication of important features, so requirements for experiments and future quantum-mechanical methods can be gauged.





Metastable $\overline{p}\alpha e$

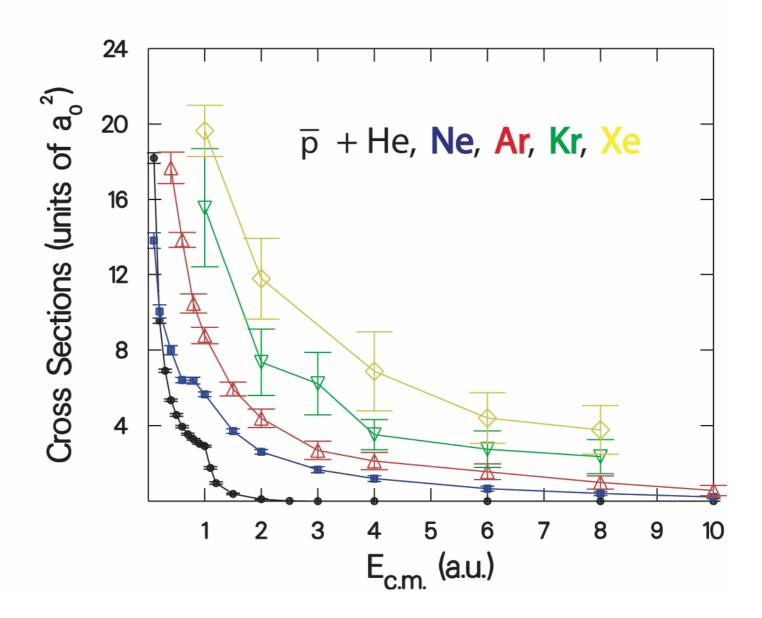


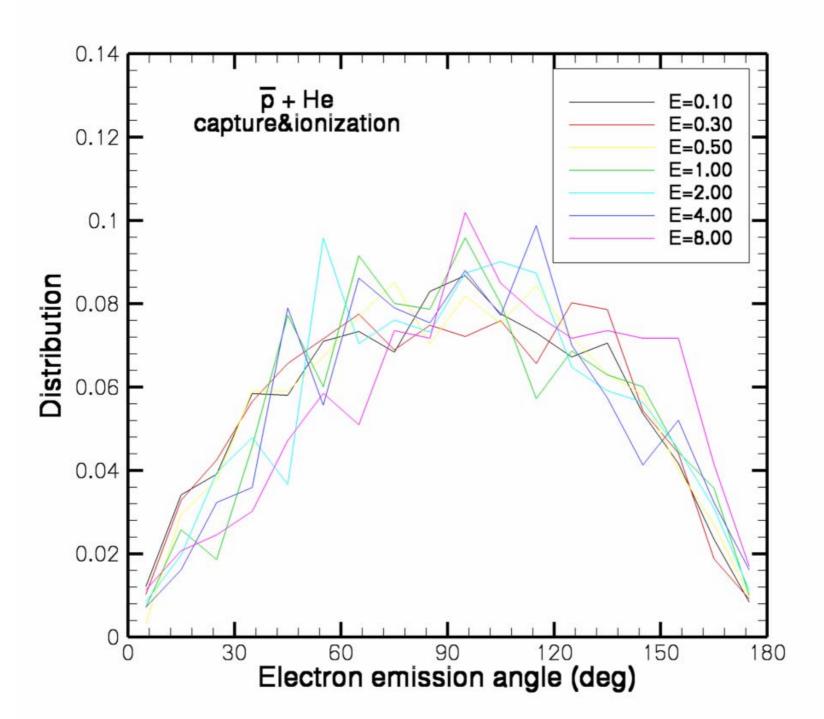
States can Auger decay if

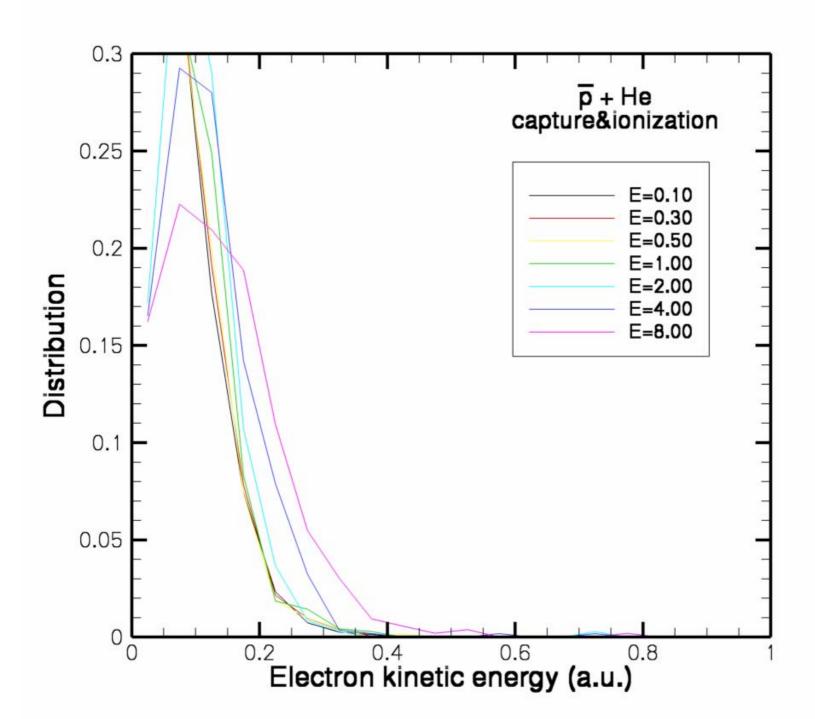
$$E_{i}^{(p)} - E_{f}^{(p)} > E_{bind}^{(e)}$$

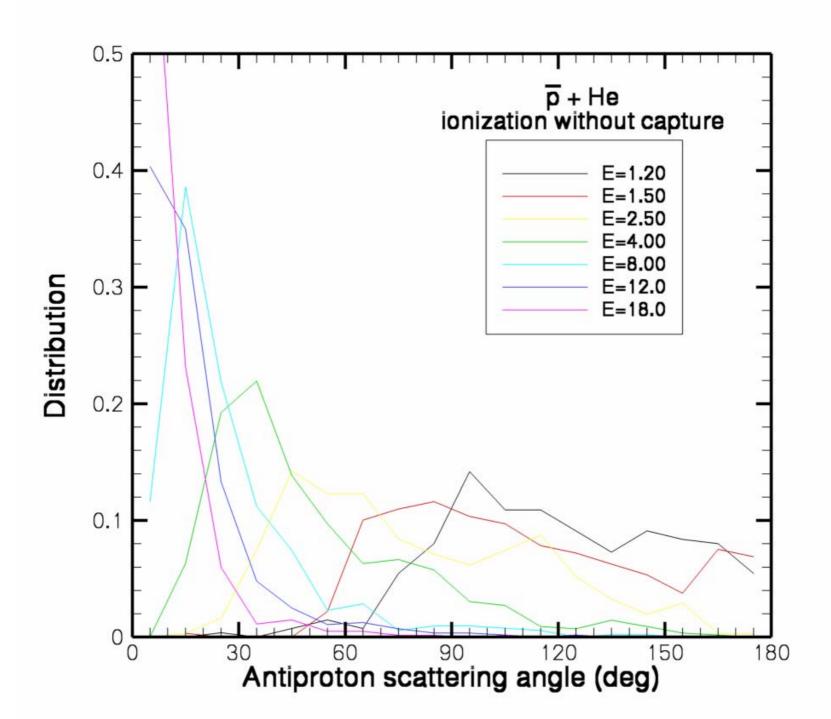
If l=n-1, then $\Delta n=1$ for dipole transitions, and this condition is satisfied only for n \lesssim 30.

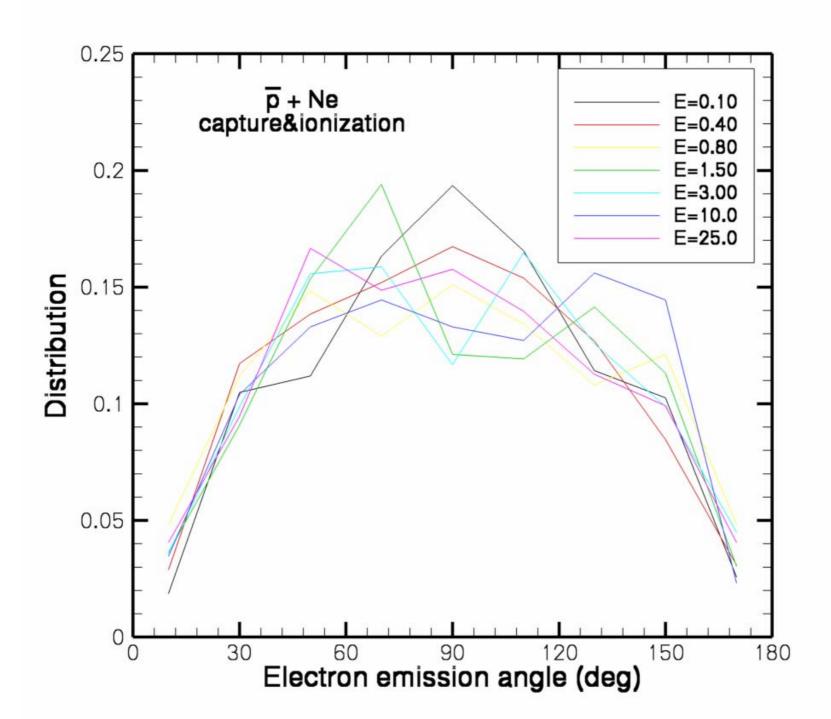
But big atoms (large *n*) may be collisionally **Stark** mixed so the selection rule doesn't apply.

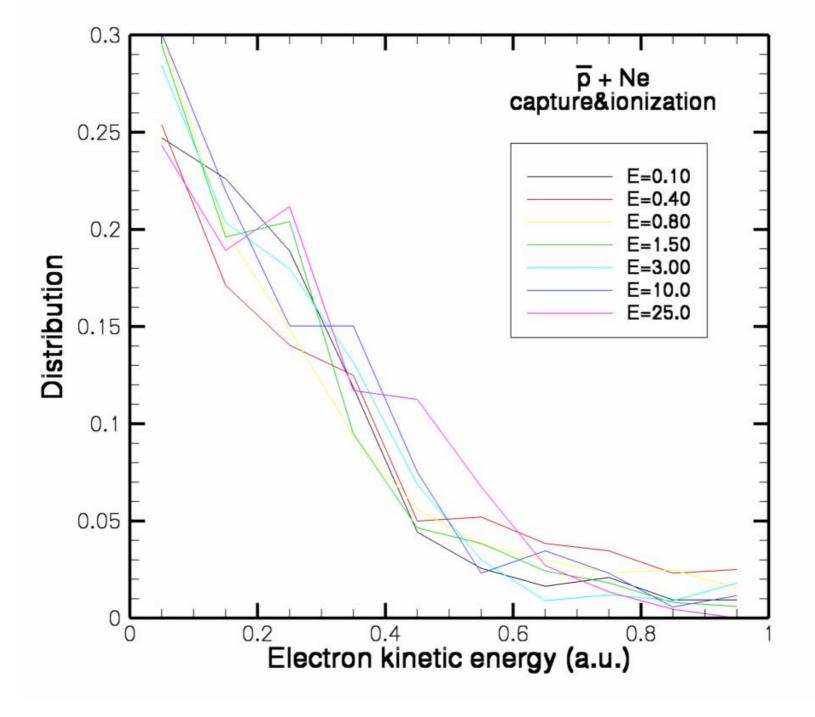


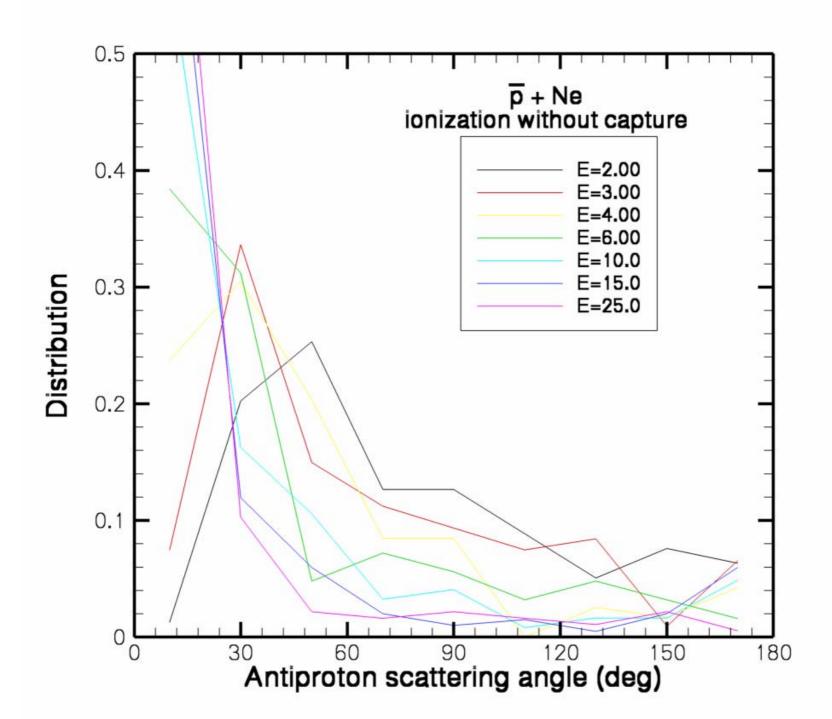




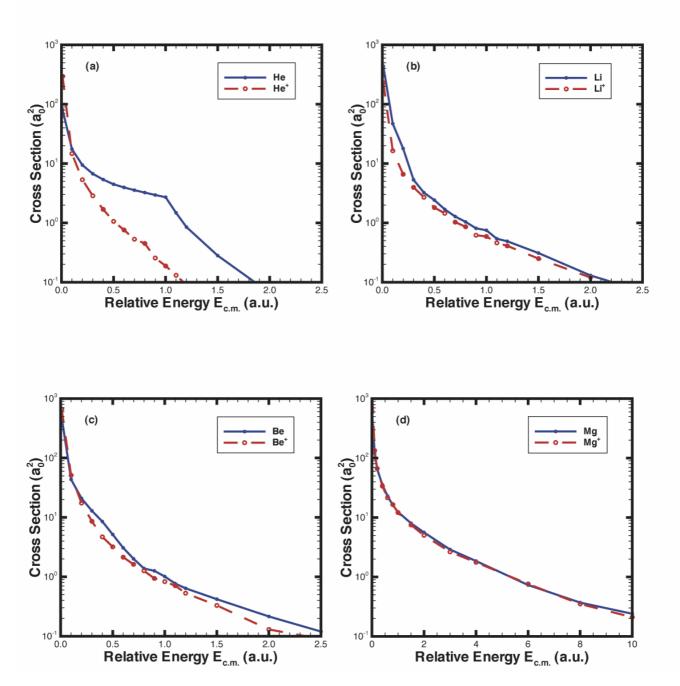




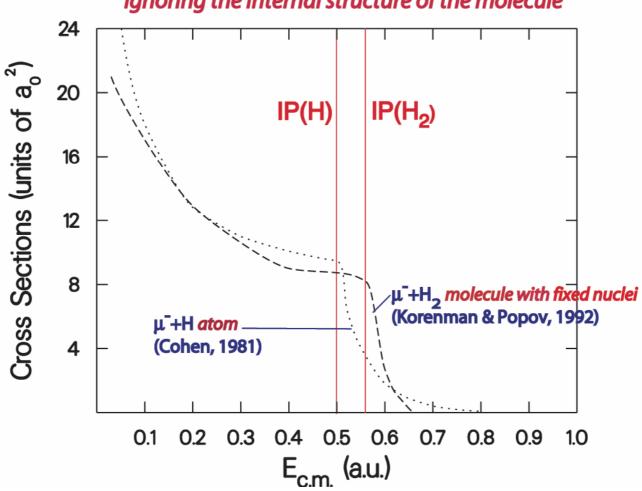




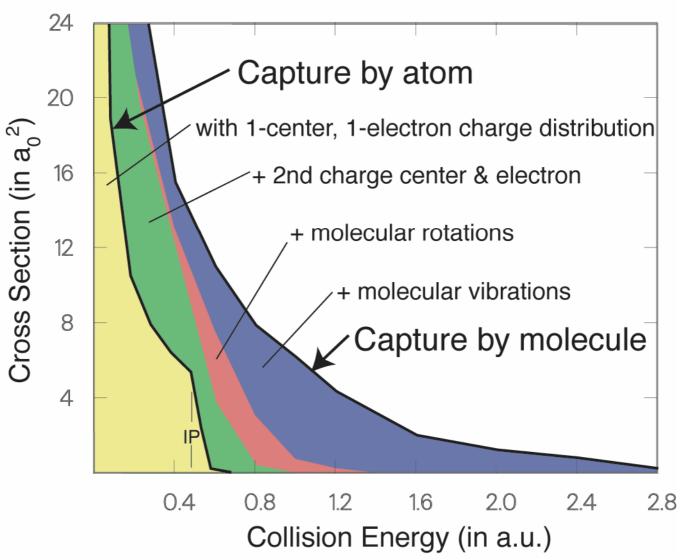
Capture of antiprotons by some (radioactive) atoms and their ions

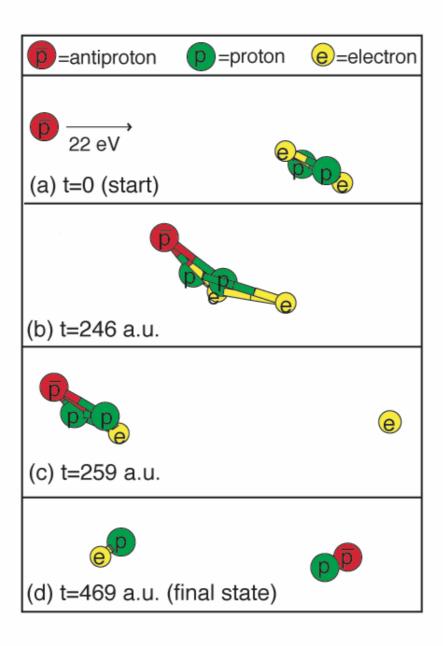


Diabatic calculations ignoring the internal structure of the molecule



$\overline{p} + H_2 \rightarrow \overline{p}p + \cdots$



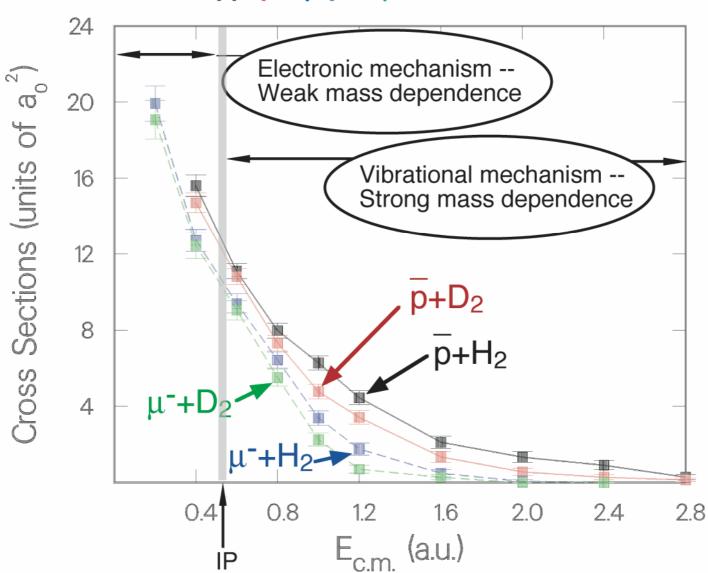


Final Arrangements in $\bar{p}+H_2$ Reaction

In capture of \bar{p} by H_2 , intermediate states such as $pp\bar{p}e^-$ or $p\bar{p}e^-$ are possible, but are predissociative and/or autoionizing. For this reason we have adopted the approach of following the trajectories long enough that the isolated $\bar{p}p$ atom can be characterized. The reactions distinguished are then

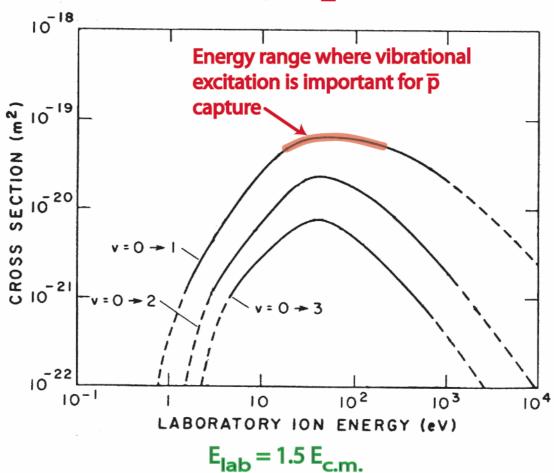
$$\bar{p} + \mathrm{H}_2 \rightarrow \begin{cases} \bar{p}p + \mathrm{H} + e^- & \sim 98\% \\ \bar{p}p + \mathrm{H}^- & \lesssim 2\% \text{ (mainly at low } E) \\ \bar{p}p + p + e^- + e^- & \lesssim 1\% \text{ (mainly at high } E) \end{cases}$$
(a) (b)

 $\overline{p}p$, $\overline{p}d$, $\mu \overline{p}$ & $\mu \overline{d}$ Formation



Comparison with *nonexotic* physics





Adapted from Figure 1 of A.V. Phelps, J. Phys. Chem. Ref. Data 19, 653 (1990).

Another comparison with *nonexotic* physics

"Lambert-Salter plot"

Collisional vibrational deactivation of polyatomic molecules at 300K

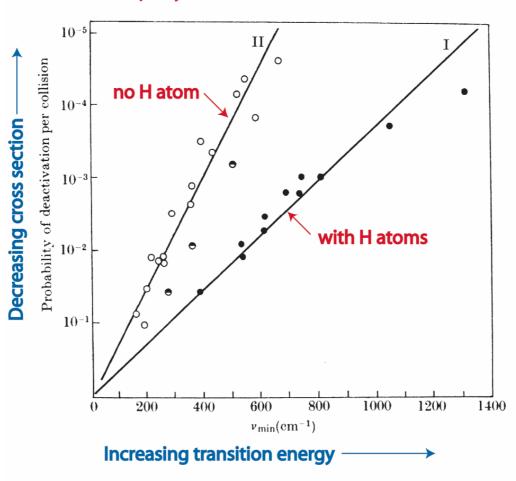


Figure 17.37 from Massey, Electronic and Ionic Impact Phenomena, vol. III

Capture of π^- in H_2+D_2 vs HD

 P_p = Probability of initial capture by p ($P_d = 1 - P_p$) W_p = Probability of final capture by p $Q = (P_p - W_p)/P_p$ = Probability of transfer $\pi^- p + d \to p + \pi^- d$

Experimental:
$$W_p^{(\text{H}_2+\text{D}_2)} = 0.405 \pm 0.010$$
 ["World Fit" $a(a-d)$] $W_p^{(\text{HD})} = 0.338 \pm 0.008^b$

Original experimental analyses assumed $P_p = P_d = 0.5$ in both H_2+D_2 and HD. With this assumption:

$$Q = (19 \pm 2)\%$$
 in H_2+D_2
 $Q = (32 \pm 2)\%$ in HD Difference was a mystery!

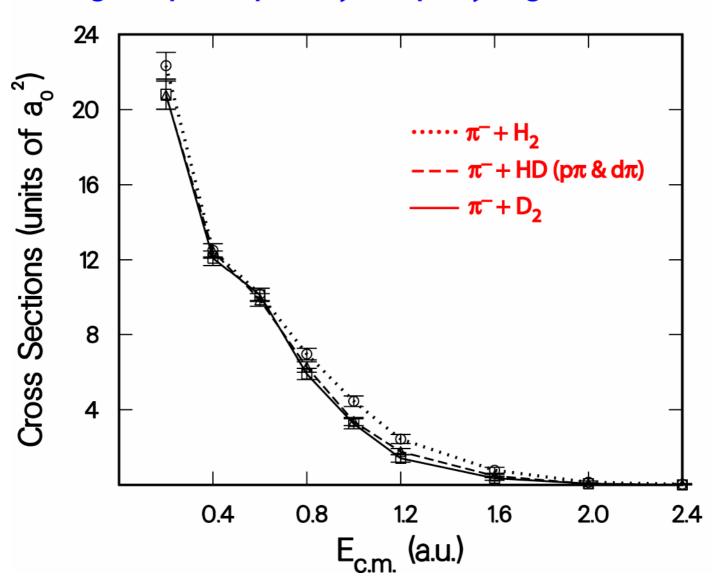
^aWeber et al. (1990)

^bAniol *et al.* (1983)

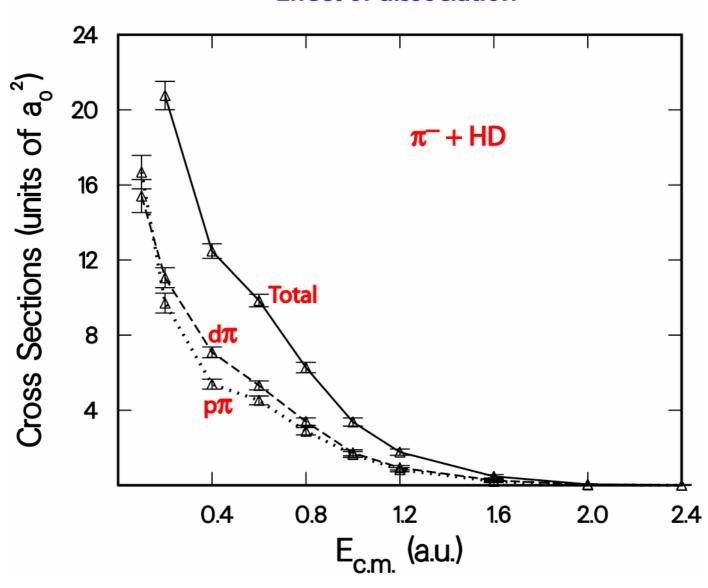
^cKravtsov et al. (1988)

^dPetrukhin & Prokoshkin (1969)

Negative pion capture by isotopic hydrogen molecules



Effect of dissociation



Capture of π^- in H_2+D_2 vs HD

 P_p = Probability of initial capture by p ($P_d = 1 - P_p$) W_p = Probability of final capture by p $Q = (P_p - W_p)/P_p$ = Probability of transfer $\pi^- p + d \to p + \pi^- d$

Experimental:
$$W_p^{(\text{H}_2+\text{D}_2)} = 0.405 \pm 0.010$$

 $W_p^{(\text{HD})} = 0.338 \pm 0.008$

Original experimental analyses assumed $P_p/P_d=1.0$ in both H_2+D_2 and HD. With this assumption:

$$Q = (19 \pm 2)\%$$
 in H_2+D_2
 $Q = (32 \pm 2)\%$ in HD Difference was a mystery!

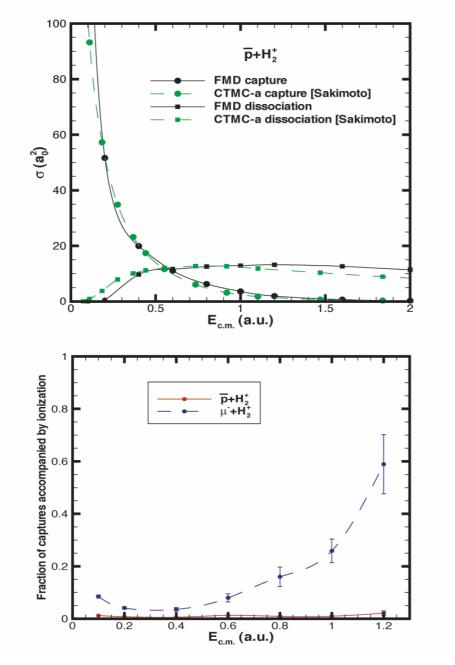
Our FMD calculations give:

$$P_p/P_d = 1.204$$
 in H₂+D₂
 $P_p/P_d = 0.875$ in HD

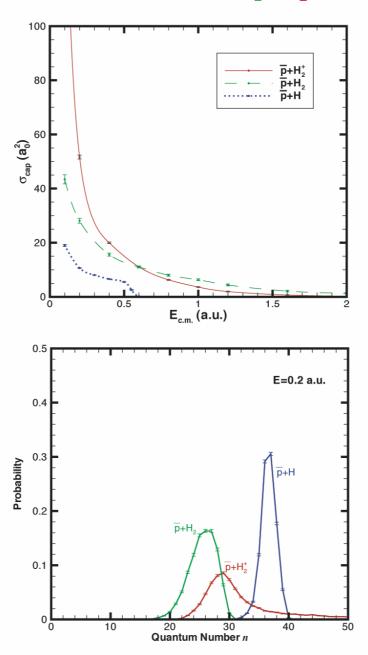
With these values of P_p/P_d , the same experimental data implies:

$$Q = (26 \pm 2)\%$$
 in H_2+D_2 Consistency is now quite satisfactory. $Q = (28 \pm 2)\%$ in HD

Nonadiabaticity & ionization in \overline{p} capture by H_2^+



Capture of \overline{p} by H, H₂ & H₂⁺



Conclusions

For capture by atoms:

- 1. The adiabatic-ionization model is valid only for the H atom. No one-electron model can be adequate for Z > 2.
- 2. The isotope effect is quite small for atoms.
- 3. The initial angular momentum distribution increases more rapidly than (2l + 1).

For capture by atomic ions:

1. Except for He⁺, capture is similar to the parent atom.

For capture by H_2 , D_2 , or HD:

- 1. Capture cross sections are larger and extend to higher collision energies for the molecular targets than for the corresponding atomic target. This effect is greater for \bar{p} than for μ^- (or π^-).
- 2. The isotope effect is large for H(or D)-containing molecules.
 - Initial capture favors the lighter nucleus.
 - Subsequent dissociation favors the heavier nucleus.
- 3. The n and l distributions are narrower and shifted to lower values relative to the distribution for the corresponding atom.

For capture by H_2^+ :

1. Capture by H₂⁺ is almost entirely due to target dissociation, not ionization.

Key features for verification in \bar{p} capture experiments

For capture by atoms and atomic ions:

- 1. H and He atoms: Rather sharp cutoff at E > I.P.
- 2. Higher-Z atoms: Multiple ionization contributes to capture.
- 3. Atomic singly charged ions: Similar to parent atom, *except* for He.

For capture by hydrogen molecules and ions:

- 1. H₂ molecule: Capture at much higher energies than for H atom.
- 2. H₂⁺ molecular ion: Capture occurs mostly without ionization.

Speculations: (based on interpretation of existing calculations)

- 1. Capture by molecules not containing H or D, e.g. O₂, may be similar to that of the constituent atoms.
- 2. Hydrides may be *qualitatively* like H_2 . (H_2O might be interesting.)

Of great basic interest:

• Comparison of cross sections and quantum-number distributions for capture of \bar{p} by H, H₂, and H₂⁺.