Lorentz and CPT Tests Involving Antiprotons

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- A. Introduction and Motivation
- B. The Standard-Model Extension (SME)
- C. Antihydrogen Spectroscopy
- D. Penning-Trap Tests
- E. Summary

A. Introduction and Motivation

Local, point-particle quantum field theories:

CPT theorem (Pauli, Lüders, Bell, '54): "Lorentz symmetry implies CPT invariance"

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Lorentz { - rotations transf. { - boosts
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CPT transf. - charge conjugation C - parity inversion P - time reversal T
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Anti-CPT theorem (Greenberg, PRL '02): "CPT violation implies Lorentz breaking"

→ { CPT TESTS } ⊂ { LORENTZ TESTS }

Why test Lorentz/CPT symmetry?

Lorentz/CPT symmetry is cornerstone of:

- present-day physics
- many candidate fundamental theories

→ Lorentz/CPT symmetry must be tested

Why look for Lorentz/CPT violation?

Nongravitational physics is well described by Standard Model (SM),

but: - phenomenological (many parameters)

- several distinct interactions

- excludes gravity

Solution: look for more fundamental theory

Candidates: string (M) theory, loop gravity, varying scalars, ...

Problem: Planck-scale measurements

(attainable energies ≪ Planck scale)

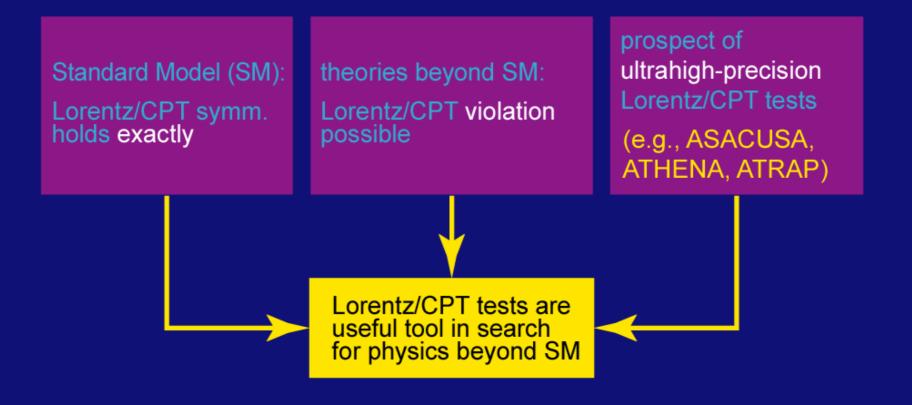
Idea: experimentally check relations that

hold exactly in Standard Model

- may be violated at fundamental level

- can be measured with ultrahigh precision

Lorentz/CPT symmetry satisfy these criteria:



Some mechanisms for Lorentz/CPT violation:

String field theory (Kostelecký et al. '90) nontrivial vacuum through spontaneous LV

Spacetime foam (Ellis et al. '98) nontrivial vacuum through virtual black holes

Nontrivial spacetime topology (Klinkhamer '00) nontrivial vacuum through compact conventional dim.

Loop quantum gravity (Alfaro et al. '00) nontrivial vacuum through choice of spin-network state

Varying scalars (Kostelecký, R.L., Perry '02) nontrivial vacuum through gradient of scalar

B. The Standard-Model Extension (SME)

Why low-energy effective theory?

Prediction of observable effects

Example: CPT ⇒ particle mass = antiparticle mass CPT ≠ particle mass ≠ antiparticle mass

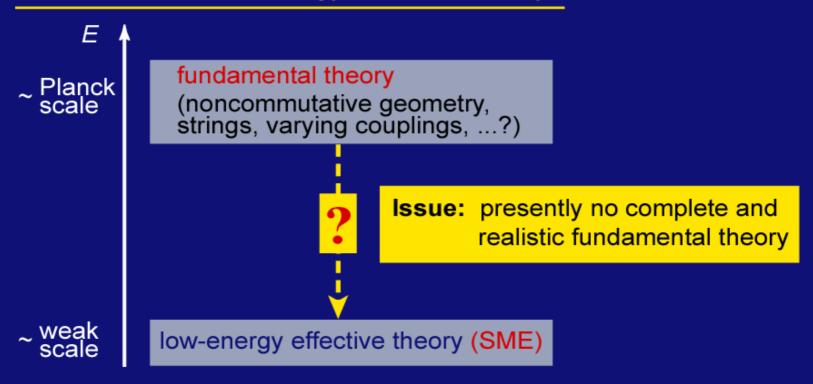
Relation between CPT-/Lorentz-breaking coefficients in different physical systems

Example: CPT-/Lorentz-breaking coefficients of quarks determine conventional, phenomenological CPT-violating parameters of kaon system

Insight into underlying theory

Example: stability / causality constrain underlying physics

How to obtain low-energy effective theory?



- Idea: examine manifestations of Lorentz/CPT violating vacuum
 - construct all possible modifications to SM

Advantage: - independent of underlying theory

- describes all low-energy effects of Lorentz violation

Construction of the SME

Definition:
$$\mathcal{L}_{\text{SME}} = \mathcal{L}_{\text{SM}} + \delta \mathcal{L}$$
 where $\delta \mathcal{L}$ contains all operators of the form (tensorial background) covariantly (SM fields)

Sample terms:

$$b^{\mu}\overline{\psi}\gamma^{5}\gamma_{\mu}\psi$$
, $(b^{\mu}\overline{\psi}\gamma^{5}\gamma_{\mu}\psi)^{2}$, $c^{\mu\nu}\overline{\psi}i\gamma_{\mu}D_{\nu}\psi$, $(k_{F})_{\mu\nu\rho\sigma}F^{\mu\nu}F^{\rho\sigma}$, ...

Remarks:

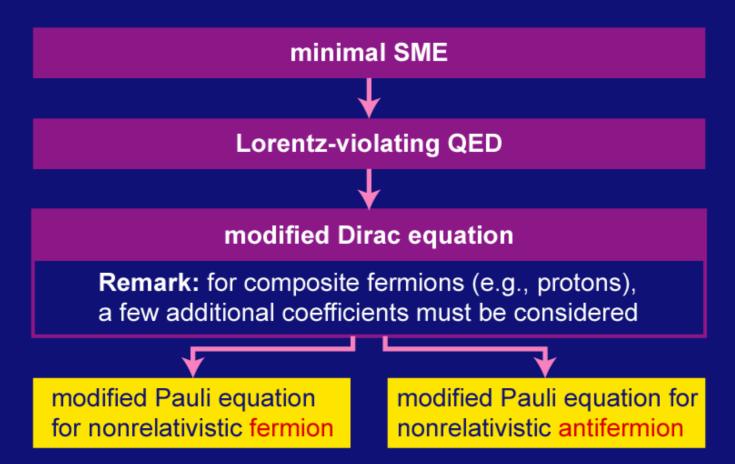
- at low E ($\ll M_{Pl}$), the background is taken as constant and the renormalizable sector of the SME dominates

minimal SME

- other features (gauge and transl. invariance, ...) can be imposed

(Colladay, Kostelecký '97; '98)

How to get SME predictions for low-energy fermions?



These modified Pauli equations are employed for Lorentz/CPT studies with cold protons/antiprotons

Sample consequences of the SME

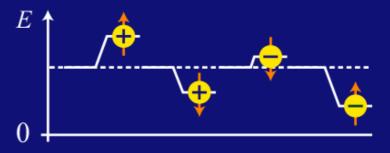
(i) modified, Lorentz-violating dispersion relations

Example:

general fermion dispersion relation (DR)

$$(p^2 - m^2)^2$$
 + CPT-/Lorentz-breaking corrections = 0

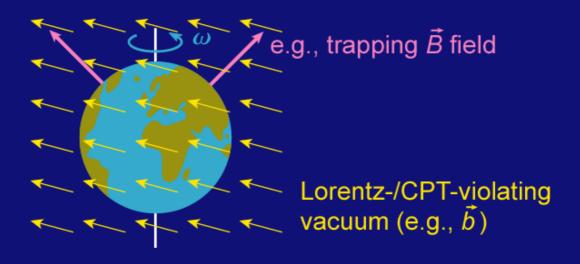
in general, 4-fold degeneracy of $E(\vec{p} = \text{const.})$ is lifted



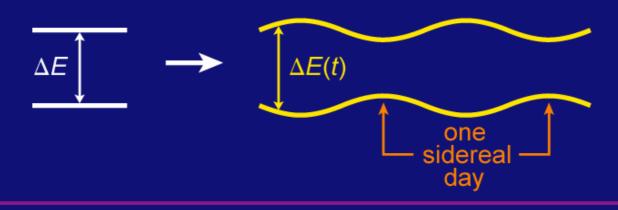
Remark:

can lead to difference in spectra for H and H (see Part C)

(ii) diurnal variations of observables



observables $\sim \vec{b} \cdot \vec{B}$ (e.g., transition frequencies) are time dependent:



Some experiments analysed within the minimal SME

Studies of neutral-meson systems

Kostelecký *et al.* '95; '96; '98; '00 OPAL Collaboration, Ackerstaff *et al.* '97 DELPHI Collaboration, Feindt *et al.* '97 KTeV Collaboration, Hsiung *et al.* '99

Tests involving photons and radiative effects

Carroll, Field, Jackiw '90 Colladay, Kostelecký '98 Jackiw, Kostelecký '99 Kostelecký, Mewes '01; '02 Kostelecký, R.L., Perry '02 Müller *et al.* '03 Lipa *et al.* '03

Penning-Trap experiments (see Part D)

Bluhm, Kostelecký, Russell '97; '98 Gabrielse *et al.* '99 Mittelman *et al.* '99 Dehmelt *et al.* '99

Hydrogen and Antihydrogen spectroscopy (see Part C)

Bluhm, Kostelecký, Russell '99 Phillips *et al.* '01

Studies of muons

Bluhm, Kostelecký, Lane '99 Hughes *et al.* '00

Clock-comparison tests

Kostelecký, Lane '99 Hunter *et al.* '99 Stoner '99 Bear et al. '00

Studies of baryogenesis

Bertolami et al. '97

Studies of neutrinos

Coleman, Glashow '99 Barger, Pakvasa, Weiler, Whisnant '00 Kostelecký, Mewes '03; '04

Kinematical studies of cosmic rays

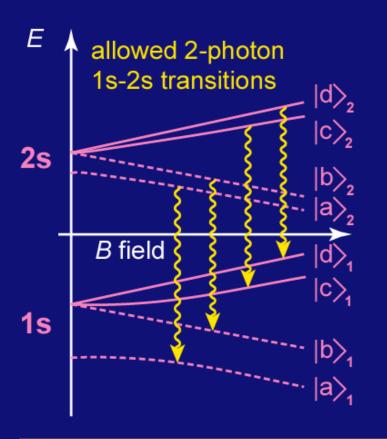
Coleman, Glashow '99 Bertolami, Carvalho '00 R.L. '03 Jankiewicz, Buniy, Kephart, Weiler '04

Tests on the ISS

Kostelecký *et al.* '02; '03 ACES PARCS RACE SUMO (OPTIS?)

C. Antihydrogen Spectroscopy

The 1s-2s transition



only the c, d states are trapped

$$|d\rangle_n = |\frac{1}{2}, \frac{1}{2}\rangle$$
 Note: no spin mixing

$$|c\rangle_n = \sin \theta_n |-\frac{1}{2}, \frac{1}{2}\rangle + \cos \theta_n |\frac{1}{2}, -\frac{1}{2}\rangle$$

with $\tan 2\theta_n \simeq \frac{51\text{mT}}{n^3 B}$

Note: θ_n , and thus spin mixing, depends on level n and field B

How are d→d and c→c transitions affected by Lorentz/CPT violation?

The $d_2 \rightarrow d_1$ transition with Lorentz/CPT violation

Leading-order energy shifts (Bluhm, Kostelecký, Russell, PRL '99)

Hydrogen (electron and proton angular momenta J and I):

$$\Delta E_{LV} = \Delta E_{e+p} + \Delta E_{e} \frac{m_{J}}{|m_{J}|} + \Delta E_{p} \frac{m_{I}}{|m_{I}|}$$

level-independent combinations of Lorentz-/CPT-violating SME coefficients

Note: both d_1 and d_2 have $m_J = 1/2$ and $m_I = 1/2$ shift is level independent

Result: no leading-order Lorentz/CPT violation in d₂ →d₁ transition

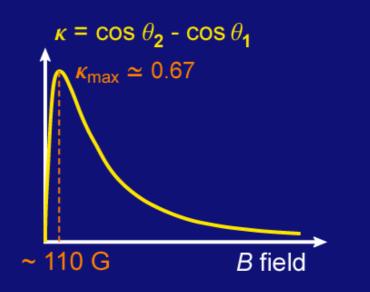
The $c_2 \rightarrow c_1$ transition with Lorentz/CPT violation

Difference between H and H transition frequencies (Bluhm, Kostelecký, Russell, PRL '99):

level-dependent spin mixing

→ unsuppressed signal

$$\Delta E_H - \Delta E_{\overline{H}} \simeq \kappa \Delta E_{\text{e+p}}$$
combination
of Lorentz-/CPT-violating
SME coefficients



Result: - leading-order Lorentz/CPT violation in c₂ →c₁ transition

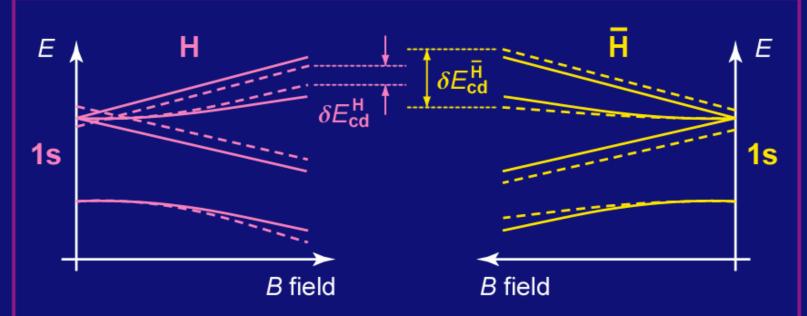
- experimental issue: effect is B-field dependent

Hyperfine Zeeman transitions within the 1s state

Difference between H and H

d1 → c1 transition frequencies

(Bluhm, Kostelecký, Russell, PRL '99; Hayano's talk tomorrow):



at field-independent transition point ($B \approx 0.65 \text{T}$):

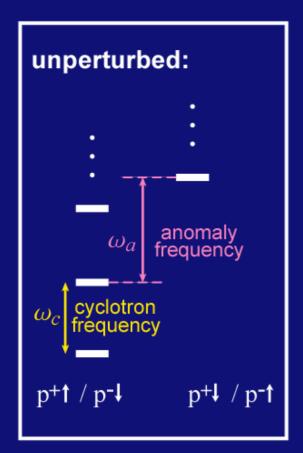
 $\delta E_{cd}^{H} - \delta E_{cd}^{H} \simeq (CPT-/Lorentz-violating SME coefficient for p)$

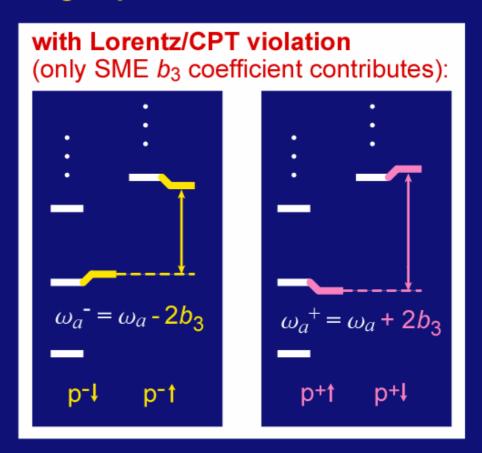
instantaneous comparison assuming 1m Hz resolution:

10⁻¹⁷ eV sensitivity to | CPT-/Lorentz-violating SME coefficient for p |

D. Penning-trap tests

Quantum-mechanical Penning-trap levels:





instantaneous comparison assuming 2 Hz resolution: 10^{-15} eV sensitivity to $|2b_3|$

E. Summary

observational tests of Lorentz/CPT symmetry are essential

- this symmetry is key ingredient in established physics
- promising tool in search for Planck-scale effects

test model for Lorentz/CPT violation is the SME

- extends our established basic physics laws to include leading-order Lorentz/CPT violation
- most general effective-field-theory test model
- has been basis for numerous Lorentz/CPT tests

cold antiprotons offer Planck-reach Lorentz/CPT tests

- unsuppressed effects in 1s-2s transition
- others in 1s hyperfine Zeeman transition
- bound parameter combinations inaccessible by other expts